

H₂ Production in the $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ Reaction in Water

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Abstract We demonstrate a method for measuring the H₂ produced in water from the $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ fission reaction. Low energy neutrons from the NIST Center for Neutron Research interact with borate-containing water in a temperature-controlled high pressure cell made from titanium. After exposure for one to several hours, the water is extracted and sparged with argon. H₂ entrained in the sparging gas is sampled with a small mass spectrometer. To determine the neutron exposure, a small amount of sodium is included in the borate solution. The water is collected and ^{24}Na activation is measured in a counting apparatus on the following day. The G-value for H₂ at room temperature is found to be (1.18 ± 0.10) molecules H₂/100eV, in good agreement with previous estimates and recent modeling calculations.

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1. Introduction

The boron-10 (^{10}B) nucleus in the chemical form of boric acid is injected into commercial nuclear reactor cooling water as a neutron "shim" to even out the neutron flux over the course of a fuel cycle.(EPRI, 2014) The $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ fission reaction usually produces a ^4He ion with kinetic energy 1.472 MeV and ^7Li ion with kinetic energy 0.840 MeV (with emission of a 0.478 MeV gamma ray); there is a 6.3 % branch of the reaction that produces a 1.776 MeV ^4He ion and a 1.015 MeV ^7Li ion and no gamma.(Barth et al., 1990) These high-LET (LET = linear energy transfer) particles quickly lose energy by ionization of the water medium, resulting in water radiolysis and production of H_2 , H_2O_2 , $\cdot\text{OH}$ radicals, $\cdot\text{H}$ atoms, and hydrated electrons in dense tracks. (Mozumder, 1999) It is straightforward to estimate the total energy deposited into the cooling water from this reaction given the nuclear cross-sections, (Brown et al., 2018; Carlson, 2011; Carlson et al., 2018) the neutron flux, and the boron concentration in the water (EPRI, 2014). At the beginning of a fuel cycle, dose from the $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ fission can amount to 30 % of the total radiation dose in the core. (Christensen, 1994; Garbett et al., 2000)

The majority of radiation dose in reactor cooling water occurs via low-LET gamma radiolysis, producing a large yield of the free radicals $\cdot\text{OH}$, $\cdot\text{H}$, and $(\text{e}^-)_{\text{aq}}$ which escape immediate recombination because of their low initial density. In stark contrast, the great preponderance of "escape" product from the high-LET ^4He and ^7Li tracks, is H_2 and H_2O_2 (in nearly equal amounts) produced by prompt recombination of the radicals in the dense tracks. (LaVerne, 2000; LaVerne, 2004) Given the importance of this source for corrosive (Lin, 2000; Macdonald, 1992; Raiman et al., 2017) H_2O_2 , it ought to be accurately included in models of the cooling water radiation chemistry (Elliot and Bartels, 2009). However a literature search indicates there exist no measurements of any product yields for this radiolysis event in high temperature water, and almost no measurements at room temperature (Barr and Schuler, 1959; Yokohata and Tsuda, 1974). H_2O_2 is very difficult to measure accurately at high temperature and pressure because it catalytically

decomposes on metal oxide surfaces and is lost to corrosion reactions (Kanjana et al., 2013; Lin et al., 1991; Satoh et al., 2004). By mass balance, the H₂ yield should be virtually identical to the H₂O₂ yield, and is far easier to measure reliably (Edwards et al., 2007; Janik et al., 2007; Sterniczuk et al., 2016). In this paper we demonstrate a methodology for measuring the H₂ yield from the $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ reaction at room temperature that can easily be extended to reactor temperatures in future work.

Unless otherwise noted, uncertainties in this paper are $k=1$, i.e., they define intervals around the central value within which the actual value is believed to lie with a level of confidence of approximately 68 percent.

2. Experimental Method

Neutron irradiations were carried out on the Beam Tube 2 (BT-2) Neutron Imaging Facility at the NIST Center for Neutron Research (NCNR). This facility has been described elsewhere (Hussey et al., 2005; J. M. LaManna, 2017). The beamline consists of a neutron beam flight tube containing a neutron beam shutter and collimator. In the present study the facility was used without any neutron imaging equipment in place and the sample cell was placed directly in the path of the ca. 6 cm diameter neutron beam as close to the exit window as possible. The irradiations were timed based on the opening and closing action of the beam shutter.

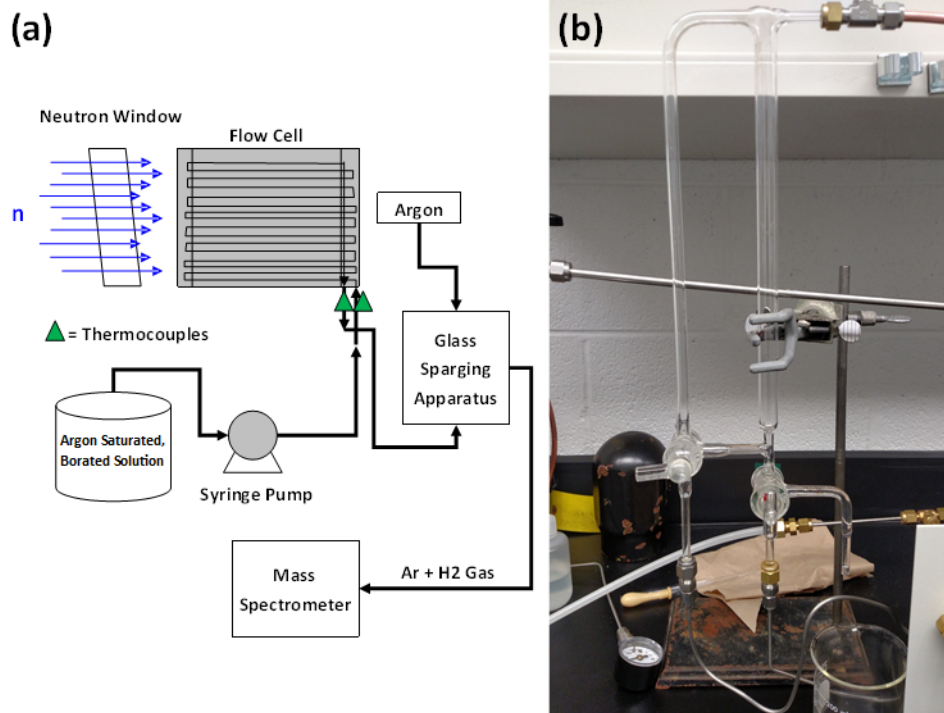


Fig. 1. A block diagram (a) of the experimental set-up for the flow system to irradiate and collect samples of borated water and an image of the glass sparging apparatus (b) used to collect the irradiated water and sparge out H_2 from the water sample.

A schematic of the apparatus used to irradiate solution, collect it, and test it for H_2 is shown in Figure 1. The irradiation cell, schematically illustrated in Figure 2, is designed for high temperature and pressures, but in this initial experiment was used at ambient conditions. It is welded together from Grade 3 titanium as a three-part "sandwich" with a close-packed honeycomb of 3.2 mm dia flow channels machined in the center section, to present the maximum liquid target to the neutron beam, but still allow fast quantitative removal of the irradiated liquid. Connecting channels are machined into the top and bottom sections of the sandwich to provide a flow-through geometry. Total liquid volume of the flow cell is ca. 90 mL. Thermocouples at the entrance and exit of the cell monitored the temperature of the water during all irradiations. All tubing used in the flow system was composed of either 316 stainless steel or Hastelloy, except for about 60 cm

of tubing connected to the inlet and outlet of the flow cell which was composed of titanium. Titanium was used for the components that would be in the direct line of neutron irradiation to reduce the total activity produced by neutron activation.

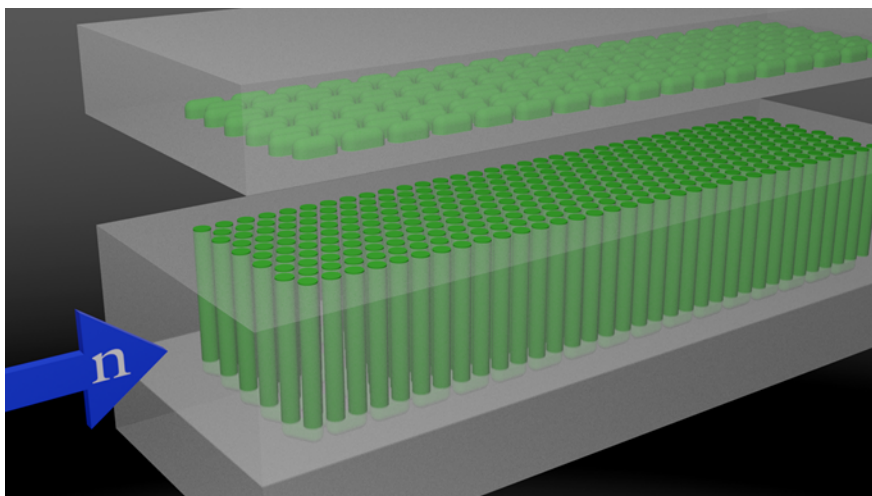


Fig 2. Schematic illustration of the three-part sandwich construction of the high pressure titanium flow cell. The close-packed flow channels (in green) contain the aqueous boric acid, and are exposed to neutrons from the end as shown.

In order to check the contribution of gamma radiolysis to the production of H_2 in the flow cell, gamma dosimetry was performed on the water-filled (no boron) flow cell using alanine pellets whose response to gamma radiation has been described previously (Sleptchonok et al., 2000). Seven alanine pellets were placed at various locations around the outside of the flow cell, specifically two near the front of the cell facing the neutron beam, one on top of the cell, one on each side of the flow cell, and two at the back of the flow cell. This setup was irradiated for 26.6 hours, after which the highest average dose recorded was (158 ± 11) Gy from the two alanine pellets positioned at the front of the cell facing the neutron beam. This result was unexpected, as the dose is of similar order of magnitude to the fission energy released by the ^{10}B , and there is virtually no

gamma incident on the cell with the neutron beam.

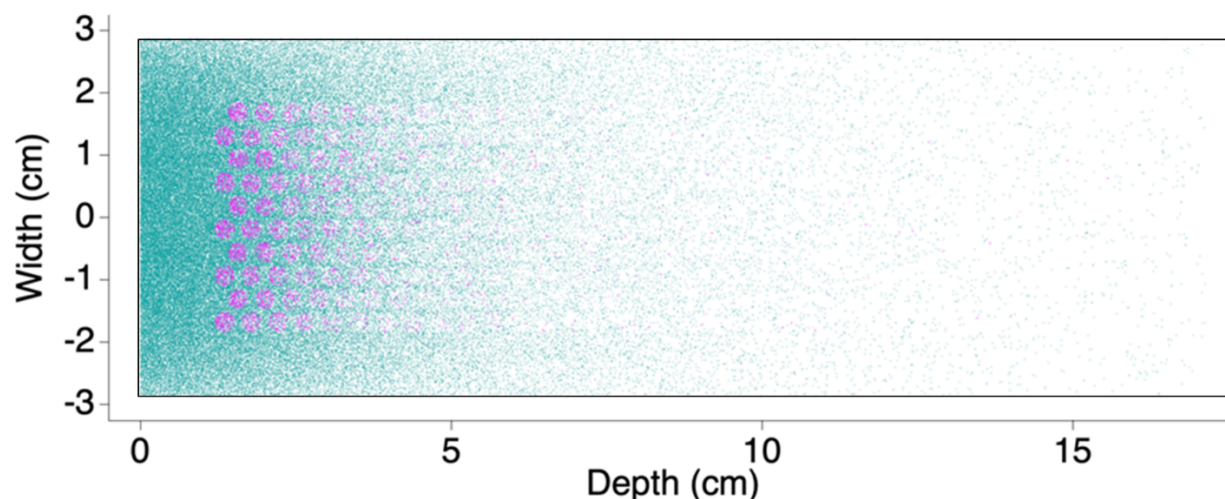


Fig 3. Plot of neutron captures in the high pressure cell by titanium (blue/green) and ^{10}B (magenta) in a center cross-section. Because there are absorptions in titanium above and below the channels there are a mixture of absorptions in those regions of the image. (See Figure 2 for a schematic drawing of the cell.)

To determine energy deposited into the boric acid solution from the products of the $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ reaction *vs.* energy released in other reactions, a geometrically-simplified model of the experimental cell and neutron beam was modeled using the Monte Carlo neutron transport code MCNP 6.2 (Werner, 2017). Figure 3 illustrates the neutron-capture events in a cross-section of the cell by either titanium (blue dots) or ^{10}B (magenta dots). The quantitative analysis relies on two F6 tallies, which sum energy deposition probabilities *vs.* energy throughout the length of a particle track. One of the F6 tallies included energy deposited in the boric acid solution from all particles, and the other included only energy deposited from the α particle and ^7Li . The ratio of integrals of the F6 tallies over energy showed the fraction of the energy deposited in the boric acid solution from the boron fission nuclei (α particle and ^7Li) is 0.71 ± 0.02 (this uncertainty is the 95 % confidence interval). Almost all of the gamma dose originates from neutron absorption and

prompt fluorescence from the major (75 %) titanium isotope ^{48}Ti . The fraction of total deposited energy from the 0.478 MeV ^{10}B gamma is approximately 1 %. Energy deposition from any absorption reactions other than titanium and boron are negligible. The MCNP calculation is used to correct the experimental result as described below.

The experimental procedure is similar to that used in a number of previous publications for both gamma and neutron radiolysis studies (Edwards et al., 2007; Janik et al., 2007; Sterniczuk et al., 2016), except that in the present case the flow is stopped during the irradiation. The borated solution was generated by mixing 40.0 g of boric acid (Sigma, BioReagent grade, ≥ 99.5 %), 0.156 g of lithium hydroxide monohydrate (Aldrich, 99.95 % trace metal basis), and 1 L of 18.2 M Ω deionized water. This solution was sparged with UHP argon for at least 30 minutes while being stirred by a magnetic stir bar. The solution was sucked out of the bottle using a Teledyne ISCO 260D syringe pump. Initially, 50 mL of solution was sent to waste out the secondary port of the syringe pump to remove any air from the pump and 200 mL of solution was flushed through the flow system to replace any solution remaining from previous runs. The total volume of the flow system is estimated to be about 120 mL from the exit port of the syringe pump to the glass sparging apparatus.

Once an irradiated solution was ready for analysis, the syringe pump was used to pump precisely 150ml of solution, including all of the irradiated volume, into a glass sparging apparatus, shown in Figure 1b. The left portion of the glass tubing allows for a continuous stream of 99.999 % argon carrier gas to bypass the solution as it is collecting in the right hand side of the apparatus. When the collection is complete, the three-way glass valve is turned to send the argon gas flow bubbling through the solution and up to the exit port. The argon and any stripped gases pass

through two Restek Molecular Sieve S-Traps in parallel to remove moisture from the gas. This gas flow then passes through a T-junction which allows sampling by an Inficon Transpector2 Compact Process Monitor Residual Gas Analyzer (RGA). Characteristic masses (typically mass 2, 28, and 32 for H_2 , N_2 , and O_2) are monitored as a function of time to determine the gas content of the water. Integration of the ion current over the peak is proportional to the aqueous concentration of that particular gas species. A typical signal for H_2 can be seen in Figure 4. Following complete analysis of the dissolved gas, the solution is sent to waste through a spigot on the glass sparging apparatus in preparation for the next run.

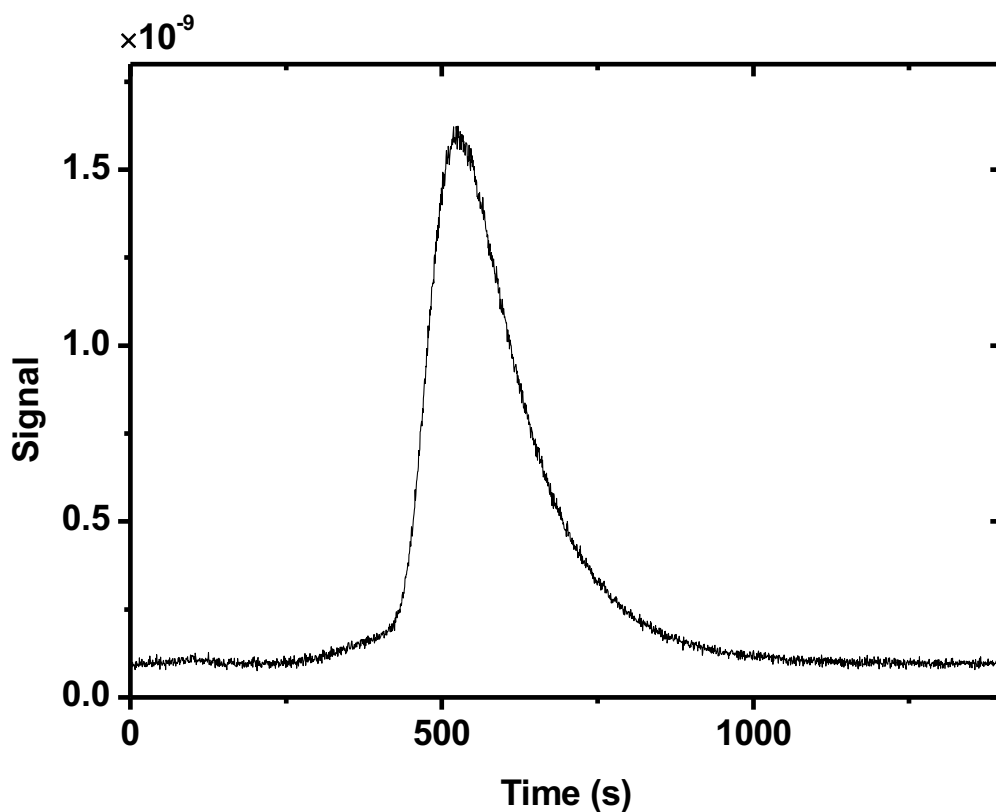


Fig. 4. A representative current signal (amps) from the mass spectrometer for $4.57 \mu\text{mol}$ of H_2 dissolved in deionized water.

To calibrate the mass spec detection of H₂, water was saturated with a gas mixture of 4.76 % H₂ in argon, obtained from Roberts Oxygen Company Inc. Different volumes of 4.76 % H₂-saturated water were collected in the sparging apparatus by controlling the flow rate and the pumping time of the syringe pump. Volumes from 25 mL to 150 mL of solution were collected and sparged. The resulting mass-2 current signals were integrated to establish a linear relationship between moles of H₂ in the sample and the area of the signal peak.

The determination of the total neutron dose was carried out via sodium activation analysis. In place of lithium hydroxide, 5.299 g of sodium carbonate (Fisher, anhydrous, Certified ACS) was dissolved with 40.0 g of boric acid in 1 L of 18.2 MΩ deionized water. This solution was sparged for at least 30 minutes with UHP argon gas while being stirred. The sodium-containing solution was irradiated for varying lengths of time from one to three hours. 150 mL of solution was collected in the glass sparging apparatus, tested for hydrogen concentration, and then collected in a plastic container. This solution was later placed in a 1 L Marinelli beaker along with 850 mL of water. Activity of the Marinelli beaker was then counted for 20 minutes on an Ortec GEM Series HPGe Coaxial Detector System gamma spectrometer (two models were used, a GEM60200-5 and a GEM40-83). These spectrometers were calibrated for activity measurements using a 1 L Marinelli activity standard provided by the National Institute of Standards and Technology. The total ²⁴Na activity was based on the weighted average value of the area under peaks located at 1368.5 and 2754.1 keV. Based on the total ²⁴Na activity produced during the neutron exposure, the average neutron flux density on the flow cell could be calculated from the following equation:

$$N_{Na-24} = \frac{N_{Na-23} \sigma_{Na-23} \Phi (1 - e^{-\lambda_{Na-24} * t})}{\lambda_{Na-24}}$$

Here $N_{\text{Na-24}}$ is the number of atoms of ^{24}Na generated in the irradiation volume, $N_{\text{Na-23}}$ is the number of initial atoms of ^{23}Na in the irradiated volume, $\sigma_{\text{Na-23}}$ is the thermal absorption cross section of ^{23}Na (0.533 barns (Brown et al., 2018)), Φ is the neutron flux density, $\lambda_{\text{Na-24}}$ is the decay rate of ^{24}Na ($1.287 \times 10^{-5} \text{ s}^{-1}$), and t is the total time of irradiation. Using the exponential decay rate, the activity ($N_{\text{Na-24}} \times \lambda_{\text{Na-24}}$) at the end of neutron exposure was precisely back-calculated from the activity measured at the time of analysis on the following day.

Results

A several-day experimental run was carried out first with several H_2 measurements vs. exposure time, and then with simultaneous H_2 and sodium activation measurements vs. exposure time. Successful calibration of the neutron flux density by the ^{24}Na activation analysis is illustrated in Figure 5, showing the good linearity between the total neutron fluence and the duration of exposure for five separate exposures.

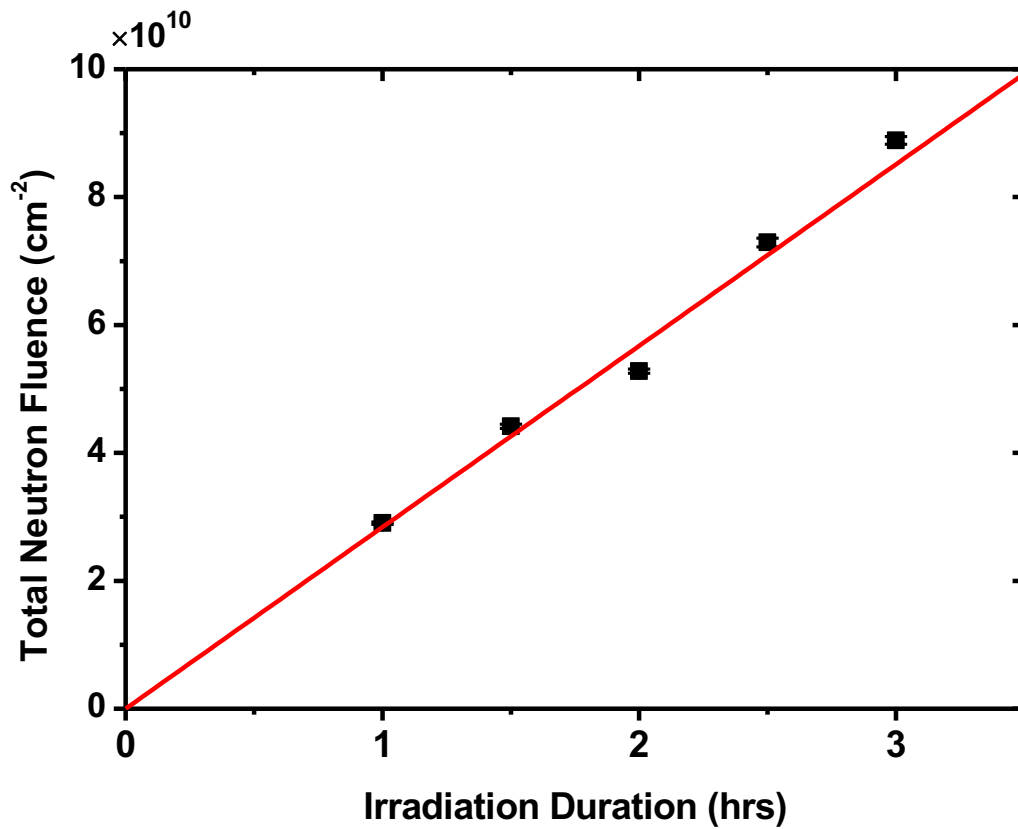


Fig. 5. Calibration of the neutron fluence versus time of irradiation. Graph shows the calculated total neutrons incident upon the cell versus irradiation time. The average neutron flux density is $(7.73 \pm 0.59) \times 10^6 \text{ cm}^{-2} \text{ sec}^{-1}$ from the slope. Statistical uncertainties are similar to the size of the plot symbols.

In order to calculate the value of $G(\text{H}_2)$ for this system, the total quantity of H_2 produced must be divided by the energy deposited in the system from the $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ reaction. The $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ reaction can proceed through two pathways: 6.3 % of the reactions release 2.791 MeV and 94 % of the reactions release 2.312 MeV as ion kinetic energy (Barth et al., 1990). The weighted average of these two energy values gives the average energy deposited in the irradiated medium per event, or 2.342 MeV. Additionally, a 0.478 MeV gamma is produced during 93.7 %

of the $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ reactions. The mean free path of these gammas is estimated to be 30 cm through water and 8 cm through titanium, so not surprisingly our simulation showed very little absorbed dose from this gamma. We therefore omit them from the G-value calculation.

The thermal capture cross section of the $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ reaction is 3869 barns according to the Evaluated Nuclear Data File neutron library (Brown et al., 2018). Due to the fact that the boric acid used in these experiments is at the natural isotopic concentration of 20 % ^{10}B , the effective thermal capture cross section is 774 barns. By using the following equation, the total energy deposited in the solution from $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ reaction can be calculated:

$$E = \Phi t \sigma_{B-\text{nat}} N_{B-\text{nat}}$$

where Φ is the average neutron flux density, $(7.73 \pm 0.59) \times 10^6 \text{ cm}^{-2} \text{ s}^{-1}$, t is the irradiation time, $\sigma_{B-\text{nat}}$ is the thermal neutron cross section of boron-natural isotopic concentration, 774 barns, and $N_{B-\text{nat}}$ is the number of boron atoms present in the irradiated boron solution prior to the irradiation. The number of hydrogen molecules sparged out of the borated solution following each irradiation can be calculated based on the calibration described previously. Figure 6 shows the hydrogen content in the borated solution versus the exposure time (upper axis) and versus total energy deposited in the solution via the $\text{B}^{10}(\text{n},\alpha)^7\text{Li}$ reaction (lower axis). The slope of the data presented in Figure 6 can be converted to a G-value for H_2 of (1.37 ± 0.05) molecules/100 eV. However, the simulations presented earlier suggest that gamma dose from ^{48}Ti neutron capture is 29 % of the total energy deposition in this experiment. Gamma $G(\text{H}_2)$ is well-known to be 0.45 ± 0.01 molecules/100eV (uncertainty is estimated from the consistent literature number) (Elliot and Bartels, 2009; Spinks and Woods, 1990; Sterniczuk et al., 2016). If we subtract this hydrogen from the measured total, the G-value for just ^{10}B neutron capture becomes (1.18 ± 0.10)

molecules/100eV.

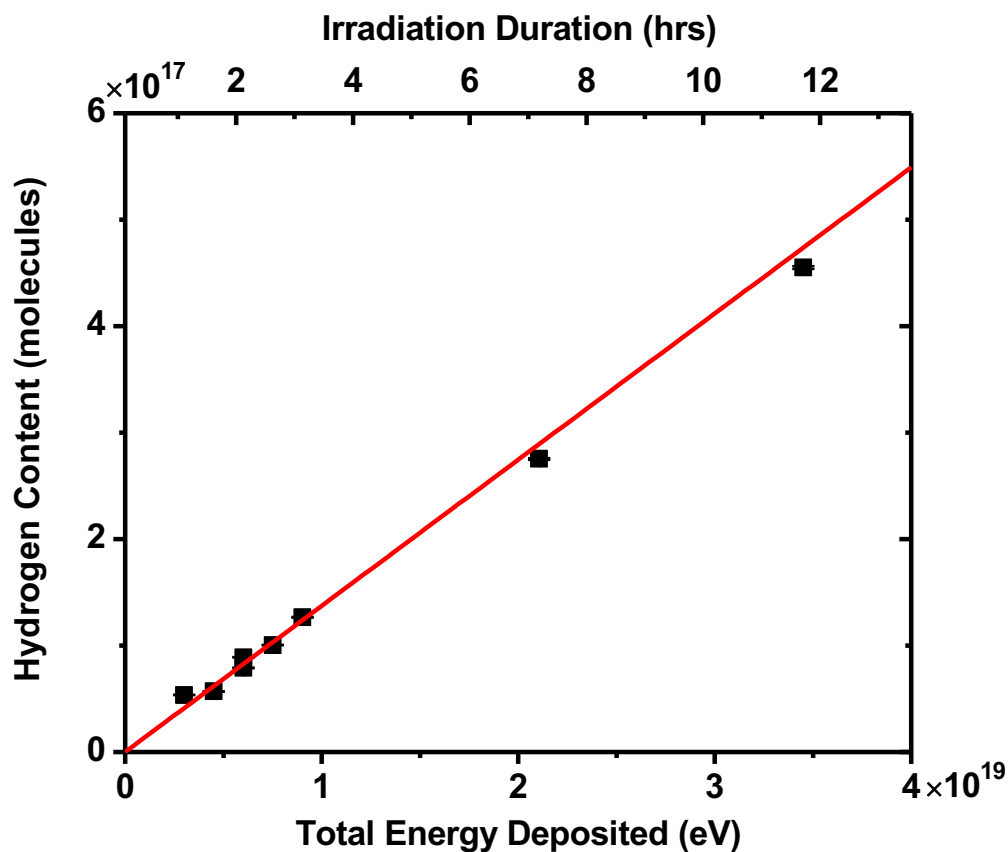


Fig. 6. Relationship between the molecules of H_2 obtained from sample versus total energy deposited in the sample via the $B^{10}(n,\alpha)^7Li$ reaction.

Discussion

As indicated in the introduction, there are very few product yield measurements for the $^{10}B(n,\alpha)^7Li$ fission reaction in aqueous solution. Barr and Schuler reported yields in the Fricke dosimeter solution (Fe^{2+} , 0.8N sulfuric acid) with and without oxygen, and Ce^{4+} reduction with and without oxygen (Barr and Schuler, 1959). The (Fe^{3+}, Ce^{3+}) products can be detected spectrophotometrically, and the four solutions allow estimates of yields $G(H_2O_2)$, $G(H)$, $G(OH)$,

and $G(-H_2O)$, where the last G-value is for decomposition of water. Much later, LaVerne and Schuler (LaVerne and Schuler, 1987) reported that the sum of oxidation yields of Fe^{3+} in 0.8N sulfuric acid from individual accelerator-produced 4He and 7Li ions gave good agreement with the $^{10}B(n,\alpha)^7Li$ fission work. To our knowledge the only $^{10}B(n,\alpha)^7Li$ radiolysis yield measurement in neutral water was a measurement of N_2 from hydrated electrons reacting with N_2O (Yokohata and Tsuda, 1974). $G(e^-)_{aq}$ was thereby estimated as 0.36 per 100 eV.

The $^{10}B(n,\alpha)^7Li$ fission product yields have been calculated recently up to 350°C with a Monte Carlo model at Sherbrooke University for both neutral water and the acidic Fricke solution (Islam et al., 2017). The assumption is made that the $^{10}B(n,\alpha)^7Li$ fission is just the sum of the individual 4He and 7Li tracks. Rather than calculate the energy-dependent differential product yields and integrate over these quantities from the initial energy down to zero (Edwards et al., 2007), the authors estimate a "representative" energy and charge for each ion, and just calculate ion radiolysis yields for that single condition at each temperature. The calculation gives yields qualitatively similar to the very limited measurements found in the literature. $G(e^-)_{aq}$ at room temperature was calculated as nearly zero at one hundred nanoseconds after the fission event, which seems to disagree with the N_2O experiment of Yokohata, and Tsuda (Yokohata and Tsuda, 1974). (This latter experiment might suffer from unexpected gamma dose much like our own.) At room temperature $G(H_2O_2)$ and $G(H_2)$ are both calculated to be ca. 1.35 molecules/100eV, in reasonable agreement with our measurement of the H_2 .

The Monte Carlo simulation (Islam et al., 2017) suggests $G(OH)$ is about 0.1 molecules/100 eV at 1 microsecond after the fission, much less than the $G(H_2)$, but not zero. More importantly, the unexpected gamma dose from ^{48}Ti neutron capture will produce $\cdot OH$ radicals with G-value of

2.6 molecules/100eV. Our initial experiment can be faulted for not including a specific $\cdot\text{OH}$ scavenger, as this radical is capable of reacting with H_2 and reducing the measured yield. However it is likely that in the presence of concentrated 0.45M borate, enough impurity is also present to react with the $\cdot\text{OH}$ instead of the H_2 . The free radical product of $\cdot\text{OH}$ reaction with nearly any impurity is unlikely to react with H_2 , so our $G(\text{H}_2)$ measurement is probably correct in spite of this experimental flaw. In fact, the 0.05M $\text{CO}_3^{=}$ counterion present in the sodium activation runs serves as an efficient $\cdot\text{OH}$ scavenger, and the H_2 yield of these runs falls on the same line in Figure 6, confirming that $\cdot\text{OH}$ production is unimportant. Any additional H_2 production from recombination of gamma-produced $(e^-)_{\text{aq}}$ and $\cdot\text{H}$ atoms is also likely prevented by impurities as well as the H_2O_2 yield (nearly equal to the H_2 yield) from ^{10}B fission.

For the moment the simulations of Jay-Gerin and coworkers remain the best $^{10}\text{B}(\text{n},\alpha)^7\text{Li}$ yield estimates for use in reactor modeling at high temperature (Islam et al., 2017). As these authors suggest, future experiments to measure the $G(\text{H}_2)$ and radical yields in high temperature water are essential to test the model predictions. Both simulations and experience with low-LET electron radiolysis (Elliot and Bartels, 2009) suggest that the free radical escape yields will tend to increase at high temperature, so a scavenger system that accounts for all of these species (Sterniczuk et al., 2016) will be essential for future experiments.

DISCLAIMER

Certain commercial equipment, instruments, or materials are identified in this paper in order to specify the experimental procedure adequately. Such identification is not intended to imply recommendation or endorsement by NIST, nor is it intended to imply that the materials or equipment identified are necessarily the best available for the purpose.

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The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript.

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ABBREVIATIONS

LET = linear energy transfer

$G(x)$ = molecules of x produced per unit of radiation energy absorbed by the system

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