# <sup>228</sup>Th breakthrough in <sup>224</sup>Ra samples: what can we know and when can we know it?

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From the medical perspective, the four alphas emitted in the Ra-224 decay represent a potent candidate for applications in targeted alpha therapy.

From the measurement perspective, the complex decay scheme makes for an interesting case; we must account for the relative detection efficiencies of all progeny and carefully account for their ingrowth.

The table shows the activity ratios predicted by the Bateman equation for a Ra-224 sample at secular equilibrium with its progeny.

The plots to the right show how progeny ingrowth proceeds over time (as calculated by the Bateman equation). In the top plot, the activities are plotted relative to the initial activity of Ra-224 at the separation time. In the bottom plot, the activities are plotted relative to the Ra-224 activity over time. In this plot we watch the relative activities converge on the values shown in the table as secular equilibrium is achieved about 6 d after separation from Th-228.

Half-lives and their (*k* = 1) uncertainties are taken from the decay data evaluation project (DDEP): <u>http://www.lnhb.fr/nuclear-data/nuclear-data-table/</u>

Uncertainties on the activity ratios are estimated by propagating nuclear data uncertainties (from DDEP) through the Bateman calculations.



The <sup>224</sup>Ra activity standard developed by NIST is based on TDCR counting. This is a selfcalibrating, "primary", liquid scintillation counting method. Measuring the number of triple- and double-coincidence events gives a measure of the counting efficiency. The source activity (A) can then be calculated from the efficiency ( $\varepsilon$ ) and count rate (N): A =  $N/\varepsilon$ 



For alpha-decays, the LS counting efficiency is assumed to be 1 (see, e.g., Fitzgerald and Forney, 2011\*\*). For the high-energy beta-decays, the LS counting efficiency is calculated using the MICELLE2 model developed by Kossert and Grau Carles (2010). For the <sup>224</sup>Ra calculations, their code was modified as described in Napoli et al., Appl. Radiat. Isotop. 155, 108933 (2020). The model provides a relationship between the TDCR and the absolute counting efficiency, which is used to calculate the sample activity.

\*\*<u>https://tsapps.nist.gov/publication/get\_pdf.cfm?pub\_id=906829</u>



The standard is based on three measurement campaigns (labeled "E2", "E3", and "E4"). In addition to TDCR, confirmatory activity measurements were performed with livetimed anticoincidence counting (LTAC), CIEMAT-NIST efficiency tracing with tritium (CNET), and  $4\pi$ - $\gamma$  measurements with an ionization chamber (AutoIC) and a well-type Nal(TI) counter. The plot shows that between-method and between-experiment agreement on the determined activity was very good.\*\*

Activities determined with high purity germanium (HPGe) detector measurements and the DDEP gamma-ray emission probability (not shown in the plot) gave a consistent 4 % bias. This bias suggests that the evaluated emission probability requires revision.

The activity standard was translated into calibration factors for radionuclide calibrators (ionization chambers) of the type found in many radiopharmacies and clinics (Napoli et al., 2020).

\*\*In this plot, the uncertainty bars for the AutoIC and NaI(TI) points are based on the statistical counting uncertainties only. These activities were calculated with benchmarked efficiency models, which carry much larger uncertainties. The TDCR, LTAC, and CNET uncertainty bars represent combined standard uncertainties (*k* = 1).

Certain commercial equipment, instruments, or materials are identified in this paper to foster understanding. Such identification does not imply recommendation by the National Institute of Standards and Technology, nor does it imply that the materials or equipment identified are necessarily the best available for the purpose.



In most experiments, we found consistent <sup>228</sup>Th breakthrough with a few Bq/MBq at the separation time. In one experiment, in April 2019, much higher levels of <sup>228</sup>Th were observed ( $f_{Th-228} = A_{Th-228}/A_{Ra-224}$ ). Dr. Roy Copping at Oak Ridge National Laboratory investigated and found that a different column type was used for that separation. The high breakthrough was considered anomalous.

The uncertainties on the impurity fractions are combined standard uncertainties (k = 1).

## Impurities matter in our measurements



 Most measurements are Liquid scintillation counting will be relatively insensitive to a most sensitive to the presence of <sup>228</sup>Th. small <sup>228</sup>Th impurity, with 1.20 corrections < 0.5 % Relative LS response f<sub>Th-228</sub> at t<sub>sep</sub> 1.15 Main contribution is delayed 0.0005 0.005 equilibrium 0.05 1.10 0.5 Becomes important for 1.05 precision comparisons or halflife determinations 1.00 5 10 15 0  $(t - t_{sep}) / d$ 

Since most measurements rely on  $\gamma$ -rays, they will be relatively insensitive to <sup>228</sup>Th (most abundant  $\gamma$ -ray, at 84 keV, is emitted in just 1.19 % of decays).

Liquid scintillation counting is sensitive to the alpha-emissions, but the very high counting efficiency of <sup>224</sup>Ra and progeny will tend to swamp the <sup>228</sup>Th signal. The plot shows the LS response of a sample with initial <sup>228</sup>Th impurity relative to a sample with no impurity. So, for example, 5 days after separation, a sample with  $f_{Th-228} = 0.05$  at  $t_{sep}$  will give about 1.1x as many LS counts as a pure Ra-224 sample. Over time, the gradual ingrowth of <sup>228</sup>Th progeny (you can think of it as 'delayed equilibration' relative to a pure <sup>224</sup>Ra source) leads to a larger effect.

'Negligible' breakthrou	gh in the literature NIST
Appl. Ruduit. Int. Vol. 39, No. 4, pp. 283–286, 1988 Int. J. Radiat. Appl. Intrum. Part A Pergamon Press pk Prieted in Great Britain	Recrived: 8 November 2017 Revised: 16 December 2017 Ascepted: 17 January 2018 DOI: 10.1002/jez.300 RESEARCHARTICLE WILEY
An Improved Generator for the Production of <sup>212</sup> Pb and <sup>212</sup> Bi from <sup>224</sup> Ra ROBERT W. ATCHER.'* ARNOLD M. FRIEDMAN'& and JOHN J. HINES' '*distion Oncology Branch, National Cancer Institute, Bethesda, Maryland and "Chemistry Division. <i>Chemistry Division Concerns Strategy</i> , Vana. <i>Chemistry Concerns</i> , 10, 6000, Vana. Weiterd 7 Conferences of the production of the strategy of the strategy of the strategy of the strategy of the thorium in the radium solution is negligible. Less than 1 ppm. Generators which have been returned to ANI. decay with the half life of 2"Ra. The yield of the generator as a function of HI	Ra-224 labeling of calcium carbonate microparticles for internal $\alpha$ -therapy: Preparation, stability, and biodistribution in miceSara Westrom <sup>1,2,3</sup> Image: A solution of the solution was estimated by the solution was estimate

Literature reports on separation of <sup>224</sup>Ra claim 'negligible' breakthrough. The Atcher (1988) paper uses apparent half-life to support this claim. The more recent Westrøm (2018) paper uses alpha spectrometry in addition to half-life.

We ask the question, "What can we know and when can we know it?"



Simulated data were generated from validated GEANT4 models of our instruments. Responses were calculated for each nuclide. Total responses at different time points were calculated by combining the individual-nuclide-responses according to the activity ratios predicted by the Bateman Equation. In this way, we could isolate the <sup>228</sup>Th contributions and discern when they might be detectable in real-world measurements.



The plot shows real data measured on a Nal(Tl) well-type counter with a "decayed" (i.e.,  $(t - t_{sep}) = 51$  d and 64 d) <sup>224</sup>Ra source. The resolution of a Nal(Tl) detector is insufficient to allow observation of the weak gamma-ray emissions from <sup>228</sup>Th. After > 10 <sup>224</sup>Ra half-lives, the spectrum appears unchanged as the contributions from <sup>228</sup>Th progeny (including, of course, <sup>224</sup>Ra and progeny) swamp the weak 84 keV gamma ray. The spectra in the figure are not normalized or decay corrected. The spectrum collected at 64 d is slightly weaker than the one collected at 51 d; the source is decaying with the <sup>228</sup>Th half-life.



The plot shows real data measured on a high purity germanium (HPGe) detector. The initial <sup>224</sup>Ra spectrum, acquired at  $(t - t_{sep}) \approx 7$  d, is shown in the black trace. After > 10 <sup>224</sup>Ra half-lives (red trace), the spectrum is preserved by <sup>228</sup>Th progeny; however, the weak gamma-ray contributions from <sup>228</sup>Th decay are resolvable.



We're defining "ppm" here, according to convention in the literature, as  $A_{Th}/A_{Ra} = f_{Th-228}$ = 10<sup>-6</sup>

Decay consistent with the <sup>224</sup>Ra half-life has been taken as the best evidence against <sup>228</sup>Th breakthrough. How useful is this metric?

Half-life measurements cannot start until the source reaches equilibrium (about 6 d after separation). Fitting datasets that include earlier points will give a biased half-life, with the magnitude of bias depending on how long you measure. Examination of the residuals of fits to datasets with early points will show obvious trends from the progeny ingrowth. It would be difficult to miss.

Fits to data acquired after equilibrium will give the correct <sup>224</sup>Ra half-life, even with  $A_{Th}/A_{Ra} > 10^{-5}$  at  $t_{sep}$ . The longer you collect data, the better your chance at detecting breakthrough.

In this plot, the DDEP-evaluated half-life is shown as a solid black line, with its (k = 1) uncertainty depicted as dashed black lines. These will appear in later plots for reference.

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Plotting what v. when

To answer "What can we know and when can we know it?", we can plot the amount of impurity that would result in a certain deviation from the <sup>224</sup>Ra half-life against the time of counting. For example, if a researcher was capable of distinguishing a deviation of just 2x the uncertainty given in the Decay Data Evaluation Project (DDEP) evaluation, and could count for 20 days after separation, a <sup>228</sup>Th impurity corresponding to 10<sup>-4</sup> at  $t_{sep}$  could be detected. These estimates are from Monte Carlo simulations of an ionization chamber, but the results would be similar for other detectors.

To give another example, if you are NOT capable of distinguishing a deviation of just 2x the DDEP uncertainty, but instead "only" capable of distinguishing a deviation at 2x the uncertainty typically achieved in a half-life determination at a national metrology institute (NMI; e.g., NIST, PTB, NPL...), then on day 20 you could detect a <sup>228</sup>Th impurity  $> 3x10^{-4}$ .



To reinforce the examples from the "What can we know and when can we know it?" plot, here are a few examples of recent half-life determinations at NIST. Also shown is the ionization chamber determination by Schrader (Appl. Radiat. Isotop. 60, 317-323 (2004)) that is the main contributor to the adopted DDEP value (the black line). The uncertainties on the determined half-life depend on the measurement conditions (initial activity, time window, etc.) and the method. It should be clear that small deviations from the <sup>224</sup>Ra half-life due to impurities would be hard to detect under any but the very best conditions.

The uncertainty bars represent combined standard uncertainties (k = 1).



The plot shows the alpha decay rate from <sup>228</sup>Th and progeny above what would be expected for a pure <sup>224</sup>Ra source. The "extra" dose is predictable if the impurity fraction, administration time, and clearance time are known.



With the methods considered here, early detection of breakthrough is not possible. If the "extra" dose from alpha particle emission from <sup>228</sup>Th and its progeny is always less than a few percent, then it may not be a concern from the medical/regulatory perspective.



## Backup slides

### NIST

## Measuring before equilibrium is tricky

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Ionization chamber response increases rapidly after separation as progeny grow in, meaning that any uncertainty on difference between  $t_{sep}$  and  $t_{meas}$  results in a very large uncertainty on the activity