Evaluation of Abramowitz functions in the right half of the complex plane

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Abstract

A numerical scheme is developed for the evaluation of Abramowitz functions J_n in the right half of the complex plane. For $n = -1, \ldots, 2$, the scheme utilizes series expansions for |z| < 1, asymptotic expansions for |z| > R with R determined by the required precision, and least squares Laurent polynomial approximations on each sub-region in the intermediate region $1 \le |z| \le R$. For n > 2, J_n is evaluated via a forward recurrence relation. The scheme achieves nearly machine precision for $n = -1, \ldots, 2$ at a cost that is competitive as compared with software packages for the evaluation of other special functions in the complex domain.

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1. Introduction

The Abramowitz functions J_n of order n, defined by

$$J_n(z) = \int_0^\infty t^n e^{-t^2 - z/t} dt, \quad n \in \mathbb{Z},$$
(1)

are frequently encountered in kinetic theory (cf., *e.g.*, [8, 17]), where the integral equations resulting from linearization of the Boltzmann equation have these functions (cf., *e.g.*, [8, 17, 26, 21]) as the kernels. The *n*-th order Abramowitz function J_n satisfies the third order ODE [1, 2]

$$zJ_{n}^{'''} - (n-1)J_{n}^{''} + 2J_{n} = 0$$
⁽²⁾

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and the recurrence relations

$$J'_{n}(z) = -J_{n-1}(z), (3)$$

$$2J_n(z) = (n-1)J_{n-2}(z) + zJ_{n-3}(z).$$
(4)

The integral representation (1) also leads to

$$J_n(\bar{z}) = \overline{J_n(z)}.$$
(5)

Research on Abramowitz functions is rather limited. In [2], about two pages of Section 27.5 are devoted to Abramowitz functions, which contain series and asymptotic expansions, originally developed in [1, 25, 38]. In [10], numerical computation of Abramowitz functions is discussed when z is a positive real number, and, in particular, it is shown that the recurrence relation for J_n is stable in both directions. In [27], a more efficient and reliable numerical algorithm using Chebyshev expansions has been developed for the evaluation of J_n (n = 0, 1, 2) when z is a positive real number.

For time-dependent or time-harmonic problems in kinetic theory, evaluation of Abramowitz functions with complex arguments is often required. However, we are not aware of any work on the evaluation of Abramowitz functions in complex domains.

In this paper, we develop an efficient and accurate numerical scheme for the evaluation of Abramowitz functions when its argument z is in the right half of the complex plane (denoted as $\overline{\mathbb{C}^+} = \{z \in \mathbb{C} | \operatorname{Re}(z) \geq 0\}$) for $n \geq -1$. We first note that Chebyshev expansions are not good representations in the complex domain since Chebyshev polynomials are orthogonal polynomials only when the argument is real. Second, when |z| is small, say, less than r for some r > 0, a series expansion can be used to evaluate $J_n(z)$ accurately with small number of terms. Third, when |z| is large, say, greater than R for some R > 0, the truncated asymptotic expansion can be used to evaluate $J_n(z)$ accurately.

We now consider the intermediate region $D = \{z \in \overline{\mathbb{C}^+} | r \leq |z| \leq R\}$, where neither the series expansion nor the asymptotic expansion can be used to achieve the required precision. Since 0 and ∞ are the only singular points of the ODE (2) satisfied by J_n , standard ODE theory [20, Chapter 16] together with the series expansion (7) shows that $J_n(z) = f_n(z) + g_n(z) \ln z$ where f and g are entire functions. Thus, J_n admits an infinite Laurent series representation in Dby theory of complex variables [5]. One may naturally ask whether $J_n(z)$ can be well approximated by a truncated Laurent series in D. It turns out that such approximation requires excessively large number of terms to achieve high accuracy. Furthermore, this global approximation is extremely ill-conditioned due to the fact that J_n behaves like an exponential function asymptotically, making its dynamic range too large to be resolved numerically with high accuracy and rendering the scheme useless.

We propose two techniques to deal with the extreme ill-conditioning associated with the global approximation of J_n in D. First, we extract out the leading factor in the asymptotic expansion of $J_n(z)$ and make a change of variable as follows:

$$J_n(z) = \sqrt{\frac{\pi}{3}} \left(\frac{\nu}{3}\right)^{n/2} e^{-\nu} U_n(\nu), \qquad \nu = 3 \left(\frac{z}{2}\right)^{2/3}.$$
 (6)

It has been shown in [1, 25] that $U_n(\nu)$ also satisfies a third order ordinary differential equation (ODE) with 0 a regular singular point and ∞ an irregular singular point. Thus, $U_n(\nu)$ is analytic for $z \in D$ and therefore can be represented by an infinite Laurent series in ν in the transformed domain. The main advantage of working with $U_n(\nu)$ instead of $J_n(z)$ is that $U_n(\nu)$ has much smaller dynamic range and thus admits more accurate and efficient approximation.

Next, we divide the intermediate region D into several sub-regions $D_i = \{z \in \mathbb{C}^+ | r_i \leq |z| \leq r_{i+1}\}$ $(i = 0, \ldots, M - 1, r_0 = r, r_M = R)$. By symmetry, we may further restrict ourself to consider the quarter-annulus domain $Q_i = \{z \in \mathbb{C} | \operatorname{Re}(z) \geq 0, \operatorname{Im}(z) \geq 0, r_i \leq |z| \leq r_{i+1}\}$ $(i = 0, \ldots, M - 1, r_0 = r, r_M = R)$. On each sub-region Q_i , we approximate $U_n(\nu)$ via a Laurent polynomial [24] in ν where the coefficients are obtained by solving a least squares problem. Here the linear system is set up by matching the function values with the values of the Laurent polynomial approximation on a set of N points on the boundary of Q_i . The least squares problem is still ill-conditioned and the conditioning becomes worse as N increases, but its solution can be used to produce very accurate approximation to the function being approximated.

Here, we would like to remark that recently least squares method has been applied to construct accurate and stable approximation for many classes of functions. In [7], it is used together with method of fundamental solutions to solve boundary value problems for the Helmholtz equation. In [15], it is used to construct rational approximation for functions on the unit circle. In [4, 3], it is shown that a wide class of functions can be approximated in an accurate and well-conditioned manner using frames and the least squares method. The least squares method is used in [16] to construct efficient and accurate sum-of-Gaussians approximations for a class of kernels in mathematical physics and in [6, 35] to construct sum-of-poles approximations for certain functions. Needless to say, the least squares problem itself has to be solved using suitable algorithms. Many such algorithms exist (see, for example, [11, 14, 18, 28, 32]).

For $n \geq 3$, we apply the recurrence relation (4) to compute $J_n(z)$. We note that the recurrence relation only needs the values of J_n for n = 0, 1, 2. Since many applications in kinetic theory require the evaluation of J_{-1} , we provide the direct evaluation of J_{-1} as well via our scheme since it is more efficient than using the recurrence relation.

Clearly, the scheme presented in this paper may be applied to the accurate evaluation of a very broad class of special functions in complex domains. Very often these special functions satisfy an ODE with a finite number of singular points. Therefore, they are analytic in complex domains excluding singular points and branch cuts. Complex analysis then ensures that Laurent series is a suitable representation to such functions in the domain. With a careful choice of the domain and suitable transformation, the least squares method becomes a reliable tool for constructing efficient, accurate and stable approximation for these functions.

The paper is organized as follows. Section 2 collects analytical results used in the construction of the algorithm. Section 3 discusses numerical algorithms for the evaluation of Abramowitz functions. Section 4 illustrates the performance and accuracy of the algorithm. The paper is concluded with a short discussion on possible extensions and applications of the work.

2. Analytical apparatus

The series expansion of J_n takes the form

$$2J_n(z) = \sum_{k=0}^{\infty} (a_k^{(n)} \ln z + b_k^{(n)}) z^k.$$
(7)

For n = 1, the coefficients can be found in [2, §27.5.4] with $a_0^{(1)} = a_1^{(1)} = 0$, $a_2^{(1)} = -1$, $b_0^{(1)} = 1$, $b_1^{(1)} = -\sqrt{\pi}$, $b_2^{(1)} = 3(1 - \gamma)/2$, and

$$a_{k}^{(1)} = -\frac{2a_{k-2}^{(1)}}{k(k-1)(k-2)}, \qquad b_{k}^{(1)} = -\frac{2b_{k-2}^{(1)} + (3k^2 - 6k + 2)a_{k}^{(1)}}{k(k-1)(k-2)}, \quad k \ge 3, (8)$$

where $\gamma \approx 0.577215664901532860606512$ is Euler's constant. For n = -1, 0, the coefficients can be obtained from term-by-term differentiation of (7), together with (3):

$$a_k^{(n)} = -(k+1)a_{k+1}^{(n+1)}, \qquad b_k^{(n)} = -(k+1)b_{k+1}^{(n+1)} - a_{k+1}^{(n+1)}, \quad k \ge 0.$$
 (9)

For n = 2, the coefficients can be obtained from term-by-term integration of (7) together with $J_2(0) = \sqrt{\pi}/4$, i.e., $a_0^{(2)} = 0$, $b_0^{(2)} = \sqrt{\pi}/2$, and

$$a_k^{(2)} = -\frac{a_{k-1}^{(1)}}{k}, \qquad b_k^{(2)} = -\frac{b_{k-1}^{(1)}}{k} + \frac{a_{k-1}^{(1)}}{k^2}, \quad k \ge 1.$$
 (10)

We have the following lemma regarding the convergence of the power series $\sum_{k=0}^{\infty} a_k^{(n)} z^k$ and $\sum_{k=0}^{\infty} b_k^{(n)} z^k$ in the series expansion (7).

Lemma 1. For $n = -1, \ldots, 2$, the power series $\sum_{k=0}^{\infty} a_k^{(n)} z^k$ and $\sum_{k=0}^{\infty} b_k^{(n)} z^k$ in (7) converge in \mathbb{C} .

Proof. For n = 1, direct calculation shows that

$$a_{2k-1} = 0, \quad a_{2k}^{(1)} = \frac{(-1)^k 2}{(2k)!(k-1)!}, \quad k > 0.$$
 (11)

Thus, the radius of convergence for $\sum_{k=0}^{\infty} a_k^{(n)} z^k$ is ∞ by the ratio test and the series converges for all complex numbers. We now split $\sum_{k=0}^{\infty} b_k^{(n)} z^k$ into the odd part and the even part:

$$\sum_{k=0}^{\infty} b_k^{(1)} z^k = z \sum_{k=0}^{\infty} b_{2k+1}^{(1)} (z^2)^k + \sum_{k=1}^{\infty} b_{2k}^{(1)} (z^2)^k.$$
(12)

For the odd part, direct calculation shows

$$b_{2k+1}^{(1)} = \frac{(-2)^k b_1^{(1)}}{(2k+1)!(2k-1)!!}.$$
(13)

Using the root test and Stirling's formula for factorials [5, p. 201], we observe that the odd part converges for all complex numbers. For the even part, we claim that

$$|b_{2k}^{(1)}| < \frac{2}{[(k-1)!]^3}, \quad k \ge 1.$$
(14)

We prove (14) by induction. First, (14) holds for k = 1 by direct calculation. Now, assume (14) holds for 2k - 2, i.e.,

$$|b_{2k-2}^{(1)}| < \frac{2}{[(k-2)!]^3}.$$
(15)

By (11), it is easy to see that

$$|a_{2k}^{(1)}| < \frac{1}{2^k [(k-1)!]^3}, \quad k > 1.$$
(16)

using the second equation in (8), we have

$$\begin{split} |b_{2k}^{(1)}| &\leq \frac{2|b_{2k-2}^{(1)}|}{2k(2k-1)(2k-2)} + \frac{3|a_{2k}^{(1)}|}{k-1} + \frac{2|a_{2k}^{(1)}|}{2k(2k-1)(2k-2)} \\ &< \frac{2|b_{2k-2}^{(1)}|}{2k(2k-1)(2k-2)} + \frac{1}{[(k-1)!]^3} \\ &< \frac{4}{2k(2k-1)(2k-2)[(k-2)!]^3} + \frac{1}{[(k-1)!]^3} \\ &< \frac{2}{[(k-1)!]^3}, \end{split}$$
(17)

where the first inequality follows from the triangle inequality, the second one follows from (16), the third one follows from the induction assumption. Thus, the even part also converges for all complex numbers by the comparison and root tests, and Stirling's formula. Finally, the convergence of the power series for n = -1, 0, 2 follows from (9), (10), (11), (13), (14), the comparison and root tests, and Stirling's formula.

Even though (7) was originally derived under the assumption that z is positive real, it makes sense for any $z \neq 0$. Furthermore, it provides a natural analytic continuation [5, p. 283] of J_n to \mathbb{C} with the branch cut along negative real axis and the principal branch for $\ln z$ chosen to be, say, $\operatorname{Im}(\ln z) \in (-\pi, \pi]$.

The asymptotic expansion of J_n is given by [2, §27.5.8]:

$$J_n(z) \sim \sqrt{\frac{\pi}{3}} \left(\frac{\nu}{3}\right)^{n/2} e^{-\nu} \left(c_0^{(n)} + \frac{c_1^{(n)}}{\nu} + \frac{c_2^{(n)}}{\nu^2} + \cdots\right), \quad z \to \infty,$$
(18)

where $\nu = 3(z/2)^{2/3}$, $c_0^{(n)} = 1$, $c_1^{(n)} = (3n^2 + 3n - 1)/12$, and

$$12(k+2)c_{k+2}^{(n)} = -(12k^2 + 36k - 3n^2 - 3n + 25)c_{k+1}^{(n)} + \frac{1}{2}(n-2k)(2k+3-n)(2k+3+2n)c_k^{(n)}, \quad k \ge 0.$$
⁽¹⁹⁾

Once again, (18) was originally derived under the assumption that z is real and positive [1, 25]. One may, however, verify that the expansion inside the parentheses on the right hand side of (18) is a formal solution to the third order ODE satisfied by U_n in (6). Furthermore, the exponential factor decays when $\arg z \in (-\frac{3\pi}{4}, \frac{3\pi}{4})$. Hence, (18) is valid for any $z \in \overline{\mathbb{C}^+}$ as $z \to \infty$. The following lemma is the theoretical foundation of our algorithm.

Lemma 2. Suppose that $D \subset \mathbb{C}$ is a closed bounded domain that does not contain the origin and the function f is analytic in D. Let $L(z) = \sum_{k=-N_1}^{N_2} c_k z^k$. Then

- (i) if $|f(z) L(z)| \le \epsilon$ for $z \in \partial D$, then $|f(z) L(z)| \le \epsilon$ for $z \in D$;
- (ii) if $|f(z) L(z)|/|f(z)| \leq \epsilon$ for $z \in \partial D$ and f has no zeros in D, then $|f(z) - L(z)| / |f(z)| \le \epsilon$ for $z \in D$.

Proof. This follows from the analyticity of L(z) on D and the maximum principle [5, p. 133].

3. Numerical Algorithms

3.1. Series and asymptotic expansions

As we have shown in Lemma 1, the coefficients $a_k^{(n)}$ and $b_k^{(n)}$ in (8)–(10) decay very rapidly and the corresponding series expansions converge for any $z \neq 0$. However, they cannot be used for numerical calculation for large |z|due to cancellation errors and increasing number of terms for achieving the desired precision. Thus, we will use the series expansions only for |z| < 1 (i.e., r = 1). In this region, both power series $\sum_{k=0}^{\infty} a_k^{(n)} z^k$ and $\sum_{k=0}^{\infty} b_k^{(n)} z^k$ converge exponentially fast and very few terms are needed to reach the desired precision. The coefficients $c_k^{(n)}$ in (19) diverge rapidly and the asymptotic expansion

(18) has to be truncated in order to be of any use. For any truncated asymptotic expansion, it is well-known that its accuracy increases as |z| increases. For a prescribed precision $\epsilon_{\rm mach}$, one needs to determine N_a — the number of terms in the truncated series, and R with |z| > R the applicable region of the truncated series. This is straightforward to determine numerically. We have found that $N_a = 18$ and R = 120 are sufficient to achieve 10^{-19} precision for J_n (n = $-1, \ldots, 2$).



Figure 1: Dynamic ranges of $J_2(z)$ and $U_2(z)$ in Q. For comparison purpose, both figures are plotted in the variable z.

3.2. Construction of the Laurent polynomial approximation for the intermediate region

We now discuss the evaluation of J_n in the intermediate region $D = \{z \in \mathbb{C}^+ | r \leq |z| \leq R\}$. First, by the conjugate property (5), we only need to discuss the evaluation of J_n in the first quadrant $Q = \{z \in \mathbb{C} | r \leq |z| \leq R, 0 \leq \arg z \leq \frac{\pi}{2}\}$. As discussed in the introduction, it is very difficult to directly approximate $J_n(z)$ in Q due to its large dynamic range. We use the transformation (6) and consider the approximation of $U_n(\nu)$ instead, U_n has a very small dynamic range. Figure 1 shows Log_{10} of $|J_2(z)|$ in Q on the left and $|U_2(z)|$ in Q on the right, where the left panel shows that the magnitude of $J_2(z)$ ranges from 10^{-19} to 10^0 , and the right panel shows that the magnitude of $U_2(z)$ ranges from 1.0 to 1.7. Other $J_n(z)$ and $U_n(z)$ exhibit similar pattern with tighter ranges for $|U_n(z)|$ (n = -1, 0, 1). Thus, we will consider the evaluation of $U_n(\nu)$ in Q.

To this end, we divide Q into several quarter-annulus domains:

$$Q_i = \{ z \in \mathbb{C} | r_i \le |z| \le r_{i+1}, 0 \le \arg z \le \frac{\pi}{2} \}, \ i = 0, \dots, M-1, \ r_0 = 1, r_M = R.$$
(20)

We will try to approximate $U_n(\nu)$ in each Q_i via a Laurent polynomial

$$U_n(\nu) \simeq L_n^{(i)}(\nu) = \sum_{k=-N_1}^{N_2} d_k^{(i)} \nu^k, \qquad z \in Q_i.$$
(21)

As noted before, $U_n(\nu)$ satisfies a third order ODE with 0 and ∞ as the only singular points [1, 25]. Thus, $U_n(\nu)$ is analytic in Q_i . By Lemma 2, in order to guarantee the accuracy of the approximation in the whole domain Q_i , it is sufficient to ensure the same accuracy is achieved on the boundary of Q_i , *i.e.*,

$$\left| U_n(\nu) - \sum_{k=-N_1}^{N_2} d_k^{(i)} \nu^k \right| \le \epsilon, \qquad z \in \partial Q_i.$$

$$(22)$$

The error-bound in (22) is achieved by solving the least squares problem

$$\mathbf{A}\mathbf{d}^{(i)} = \mathbf{f},\tag{23}$$

where

$$A_{jk} = \nu_j^k, \qquad f_j = U_n(\nu_j), \quad j = 1, \dots, 4N_b,$$
 (24)

where $\nu_j = 3(z_j/2)^{2/3}$, and z_j are chosen to be the images of Gauss-Legendre nodes on each segment of ∂Q_i , N_b is chosen to ensure that the error of approximation of $U_n(\nu)$ by the corresponding Legendre polynomial interpolation on each segment of ∂Q_i is bounded by ϵ . The right hand side **f** in (23) is computed *via* symbolic software system MATHEMATICA to at least 50 digits. In other words, we do not use the actual analytic Laurent series to approximate U_n on each quarter-annulus Q_i . Instead, a numerical procedure is applied to find much more efficient "modified" Laurent series for approximating U_n on each Q_i .

The linear system (23) is ill-conditioned. However, since we always use $\mathbf{d}^{(i)}$ in the Laurent polynomial approximation to evaluate U_n , we obtain (by the maximum principle) high accuracy in function evaluation in the entire subregion as long as the residual of the least squares problem (23) is small.

The least squares solver also reveals the numerical rank of \mathbf{A} , which is used to obtain the optimal value of $N_{\rm T} = N_2 - N_1 + 1$, the total number of terms in the Laurent polynomial approximation. It is then straightforward to use a simple search to find the value for N_1 , which completes the algorithm for finding a nearly optimal and highly accurate Laurent polynomial approximation for U_n in Q_i .

Remark 1. We would like to emphasize that the Laurent polynomial approximation may not be unique, but this non-uniqueness has no effect on the accuracy of the approximation.

Remark 2. We have computed the integrals

$$I_n = \int_{\partial Q} \frac{J'_n(z)}{J_n(z)} dz = -\int_{\partial Q} \frac{J_{n-1}(z)}{J_n(z)} dz \tag{25}$$

for n = -1, ..., 2 and found numerically that they are all close to zero. By the argument principle [5, p. 152], we have

$$I_n = 2\pi i (Z_n - P_n), \tag{26}$$

where Z_n and P_n denote respectively the number of zeros and poles of $J_n(z)$ inside ∂Q . Since $J_n(z)$ is analytic in Q, it has no poles in Q, i.e., $P_n = 0$. Thus, the fact that I_n is very close to zero shows that $Z_n = 0$, that is, J_n has no zeros in Q. Further numerical investigation shows that functions $|U_n(\nu)|$ $(n = -1, \ldots, 2)$ range from 0.95 to 1.7 on ∂Q . Combining these two facts, we conclude that the absolute error bound on the approximation of U_n gives roughly the same relative error bound.

3.3. Evaluation of J_n for $n = -1, \ldots, 2$

Once the coefficients of Laurent polynomial approximation for each subregion are obtained and stored, the evaluation of $J_n(z)$ is straightforward. That is, we first compute |z| to decide on which region the point lies, then use the proper representation to evaluate $J_n(z)$ accordingly. We summarize the algorithm for calculating $J_n(z)$ for $z \in \overline{\mathbb{C}^+}$, $n = -1, \ldots, 2$ in Algorithm 1.

Algorithm 1 Evaluation of $J_n(z)$ for $z \in \overline{\mathbb{C}^+}$

procedure ABRAM(z, f)

 \triangleright Input parameter: z - the complex number for which the Abramowitz function J_n is to be evaluated.

 $\begin{array}{ll} \triangleright \mbox{ Output parameter: } f \mbox{ - the value of Abramowitz function } J_n(z). \\ \mbox{ assert } \operatorname{Re}(z) \geq 0. \\ \mbox{ if } |z| \leq 1 \mbox{ then } & \triangleright z \mbox{ is in the series expansion region.} \\ \mbox{ Use the series expansion (7) to evaluate } f = J_n(z). \\ \mbox{ else if } |z| \geq 120 \mbox{ then } & \triangleright z \mbox{ is in the asymptotic region.} \\ \end{array}$

Set $\nu = 3(z/2)^{2/3}$.

Use the asymptotic expansion (18) to compute $U_n(\nu)$.

Set $f = \sqrt{\frac{\pi}{3}} \left(\frac{\nu}{3}\right)^{n/2} e^{-\nu} U_n(\nu).$

 $\triangleright z$ is in the intermediate region.

Set $\nu = 3(z/2)^{2/3}$.

Use a precomputed Laurent polynomial approximation (21) to compute $U_n(\nu)$.

Set $f = \sqrt{\frac{\pi}{3}} \left(\frac{\nu}{3}\right)^{n/2} e^{-\nu} U_n(\nu)$. end if

end procedure

else

Remark 3. All these expansions can be converted into a polynomial of certain transformed variable. We use Horner's method [23, §4.6.4] to evaluate the polynomial in the optimal number of arithmetic operations.

Remark 4. The accuracy of $J_n(z)$ deteriorates as |z| increases since the condition number of evaluating the exponential function $e^{-\nu}$ is $|\nu|$. This is unavoidable in any numerical scheme as the phenomenon is related to physical ill-conditioning of evaluating $J_n(z)$ for the argument with large magnitude.

3.4. Evaluation of J_n for n > 2

In [10], it is shown that (4) is stable in both directions when z is a positive real number. We have implemented the forward recurrence to evaluate $J_n(z)$ for n > 2. We have not observed any numerical instability during our numerical tests for $z \in \overline{\mathbb{C}^+}$.

4. Numerical results

We have implemented the algorithms in Section 3 and the code is available at https://github.com/zgimbutas/abramowitz. Numerical experiments were performed on a desktop with a 3.10GHz Intel(R) Xeon(R) CPU.

For the series expansion (7), a straightforward calculation shows that 18 terms in $\sum b_k^{(n)} z^k$ and 9 nonzero terms in $\sum a_k^{(n)} z^k$ are needed to reach 10^{-19} precision for J_n (n = -1, ..., 2). For the asymptotic expansion (18), we find that it is sufficient to choose $N_a = 18$, R = 120 for 10^{-19} precision. All coefficients are precomputed with 50 digit precision.

For the intermediate region, we divide |z| on [1, 120] into three subintervals [1,3], [3,15], [15,120] and Q into Q_1 , Q_2 , Q_3 , respectively. We use IEEE binary128 precision to carry out the precomputation step and solve the least squares problem with 10^{-20} threshold for the residual. We have found that for Q_1 we need $N_2 = 11$, $N_T = 30$ for J_0 and J_1 , $N_2 = 10$, $N_T = 32$ for J_{-1} , and $N_2 = 11$, $N_T = 32$ for J_2 . For all four functions J_n (n = -1, 0, 1, 2), we need $N_2 = 0$, $N_T = 30$ for Q_2 and $N_2 = 0$, $N_T = 20$ for Q_3 . The coefficients of Laurent polynomial approximations for J_n (n = -1, 0, 1, 2) on Q_i (i = 1, 2, 3)are listed in Tables B.4–B.15 in Appendix B.

Remark 5. The coefficients in Tables B.4–B.15 for Q_2 and Q_3 do not have small norms. However, for Q_2 , $\left|\frac{1}{\nu}\right| \leq \frac{1}{3(3/2)^{(2/3)}} = 0.254...$; and for Q_3 , $\left|\frac{1}{\nu}\right| \leq \frac{1}{3(15/2)^{(2/3)}} \approx 0.087$. It is easy to see that terms $c_j \left(\frac{1}{\nu}\right)^j$ decrease as j increases. Alternatively, we could consider the Laurent series of the form $\sum \tilde{c}_j \left(\frac{\nu_i}{\nu}\right)^j$ with $\nu_i = 3(r_i/2)^{(2/3)}$ (r_i is the lower bound for |z| in Q_i). Then the coefficient vector \tilde{c} will have small norm, as required in [7, 4]. However, this corresponds to the column scaling in the least squares matrix and almost all methods for solving the least squares problems do column normalization. Thus, it has no effect on the accuracy of the solution and stability of the algorithm.

Remark 6. The partition of the sub-regions is by no means optimal or unique. There is an obvious trade-off between the number of sub-regions and the number of terms in the Laurent polynomial approximation. For example, one may use a finer partition for the regions closer to the origin. We have tried to divide the intermediate region into 14 regions with $Q_i = \{z \in \overline{\mathbb{C}^+} | (\sqrt{2})^{i-1} \le |z| \le (\sqrt{2})^i \}$ (i = 1, ..., 14), and we observe that only 20 terms are needed for all regions. However, our numerical experiments indicate that the partition has very mild effect on the overall performance (i.e., speed and accuracy) of the algorithm.

4.1. Accuracy check

We first check the accuracy of Algorithm 1. The reference function values are calculated via MATHEMATICA to at least 50 digit accuracy. The error is measured in terms of maximum relative error, i.e.,

$$E = \max_{i} \frac{|\hat{J}_n(z_i) - \tilde{J}_n(z_i)|}{|\tilde{J}_n(z_i)},$$

where $\tilde{J}_n(z_i) = e^{\nu_i} J_n(z_i)$ ($\nu_i = 3(z_i/2)^{2/3}$) is the reference value of the scaled Abramowitz function computed via MATHEMATICA, and $\hat{J}_n(z_i)$ is the value computed via our algorithm. The points z_i are sampled randomly with uniform distribution in both its magnitude and angle in $\overline{\mathbb{C}^+}$. Table 1 lists the errors for evaluating \tilde{J}_n (n = -1, 0, 1, 2) in various regions, where we observe that the errors are within $10\epsilon_{\text{mach}}$ with the machine epsilon $\epsilon_{\text{mach}} \approx 2.22 \times 10^{-16}$ for IEEE double precision. In general, the errors in the first intermediate region Q_1 are slightly bigger due to mild cancellation errors.

Table 1: The relative L^{∞} error of Algorithm 1 over 100,000 uniformly distributed random points in $\overline{\mathbb{C}^+}$. The reference value is computed via MATHEMATICA to at least 50 digit accuracy. S denotes the series expansion region and A denotes the asymptotic expansion region.

	S	Q_1	Q_2	Q_3	А
J_{-1}	1.5×10^{-15}	$2.1{\times}10^{-15}$	4.4×10^{-16}	6.4×10^{-16}	8.6×10^{-16}
J_0	1.3×10^{-15}	2.4×10^{-15}	2.2×10^{-16}	2.2×10^{-16}	2.2×10^{-16}
J_1	1.1×10^{-15}	2.4×10^{-15}	4.7×10^{-16}	6.0×10^{-16}	8.0×10^{-16}
J_2	1.2×10^{-15}	2.9×10^{-15}	$5.6 imes 10^{-16}$	8.4×10^{-16}	1.2×10^{-15}

For n > 2, extensive numerical experiments indicate that the forward recurrence relation (4) is stable for evaluating J_n in $\overline{\mathbb{C}^+}$. The relative errors are shown in Table 2 for a typical run.

Table 2: The maximum relative error for evaluating J_{100} using the forward recurrence relation (4) over 100,000 uniformly distributed random points in the domain $\{z \in \mathbb{C} | \operatorname{Re}(z) \ge 0, 0 < |z| < 1000\}$. The reference values are calculated using MATHEMATICA with 240-digit precision arithmetic.

S	Q_1	Q_2	Q_3	А
1.3×10^{-15}	$2.9{\times}10^{-15}$	$1.3{\times}10^{-15}$	2.0×10^{-15}	$3.7{\times}10^{-15}$

4.2. Timing results

Since all three representations (i.e., Laurent polynomials, series and asymptotic expansions) mainly involve polynomials of degree less than 30, the algorithm takes about constant time per function evaluation in $\overline{\mathbb{C}^+}$. We have tested the CPU time of Algorithm 1 for evaluating $\tilde{J}_n(z)$ and compared it with that of evaluating the complex error function $\operatorname{erf}(z) = \frac{2}{\sqrt{\pi}} \int_0^z e^{-t^2} dt$. The complex error function is a well studied special function that has received much attention in the community of scientific computing. See, for example, [9, 12, 13, 19, 29, 30, 31, 33, 34, 36, 37]. Here we use the well-regarded Faddeeva package [22] to evaluate $\operatorname{erf}(z)$.

The results are shown in Table 3. First, we note that erf(z) is an entire function which is somewhat simpler than the Abramowitz functions and the

Faddeeva package guarantees about 10^{-13} accuracy. Second, the numbers of terms in all three representations in our algorithm are chosen so that 10^{-19} precision may be achieved if the calculation were carried out in 80-bit floating-point arithmetic (it achieves about 10^{-15} accuracy in double precision arithmetic as shown in Table 1).

In the asymptotic region, our algorithm is slightly faster than the numbers shown in Table 3, while the Faddeeva package is faster by a factor of about 3. However, the efficiency in the asymptotic region (i.e., the asymptotic expansion) heavily depends on the properties of the given special functions and is thus independent of the algorithm for other regions. Combining all these factors, we may conclude that our algorithm is competitive with the highly optimized Faddeeva package.

Table 3: The total CPU time in seconds for evaluating $J_n(z)$ using Algorithm 1 and the error function erf(z) over 1,000,000 uniformly distributed random points in $0 \leq \text{Re}(z) \leq 10$, $0 \leq \text{Im}(z) \leq 10$.

	$J_{-1}(z)$	$J_0(z)$	$J_1(z)$	$J_2(z)$	$\operatorname{erf}(\mathbf{z})$
T	0.44	0.41	0.44	0.41	0.34

5. Conclusions and further discussions

We have designed an efficient and accurate algorithm for the evaluation of Abramowitz functions J_n in the right half of the complex plane. Some useful observations in the design of the algorithm are applicable for evaluating many other special functions in the complex domain. First, it is better to pull out the leading asymptotic factor from the given function when |z| is large. Second, the maximum principle reduces the dimensionality of the approximation problem by one. Third, the least squares scheme is generally a reliable and accurate method to find an approximation of a prescribed form. That is, analytical representations should be used with caution even if they are available, as they often lead to large cancellation error or very inefficient approximations or both.

Finally, though we have used Laurent polynomials for approximating Abramowitz functions in the intermediate region, there are many other representations for function approximations. This includes truncated series expansion, rational functions (see, for example, [15]), etc. We have actually tested the truncated series expansion in the sub-region (i.e., Q_1) closest to the origin for J_n . Our numerical experiments indicate that the performance is about the same as the one presented in this paper.

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Appendix A. Zeros of $J_n(z)$

We have used NINTEGRATE in MATHEMATICA to evaluate I_n defined in (25). When WORKINGPRECISION is set to 100, $|I_n|$ are about 10^{-59} for n = -1, 0, 1, 2. When it is set to 200, the values of $|I_n|$ decrease to 10^{-160} . By the argument principle, I_n can only take integral multiples of $2\pi i$. Thus, the numerical calculation clearly shows that J_n (n = -1, 0, 1, 2) have no zeros in the intermediate region Q. Analytically, we can only show that J_n has no zeros in the sector $|\arg(z)| \leq \frac{\pi}{4}$. The proof is presented below.

Lemma 3. If $z_0 \in \mathbb{C}$ is a zero of $J_n(z)$, then also \overline{z}_0 .

Proof. This simply follows from the conjugate property (5).

Lemma 4. Suppose that $n \ge 0$. Then $J_n(z)$ has no zero in the sector $|\arg z| \le \frac{\pi}{4}$.

Proof. Let $z_0 = x_0 + iy_0 \in \overline{\mathbb{C}^+}$ be a zero of $J_n(z)$. Then by Lemma 3, \overline{z}_0 is also a zero of $J_n(z)$. Consider functions $f(t) = J_n(z_0t)$ and $g(t) = J_n(\overline{z}_0t)$. Then f(1) = g(1) = 0, and f, g and their derivatives decay exponentially fast to 0 as $t \to \infty$ by the asymptotic expansion (18).

The differential equation (2) implies that

$$tf'''(t) - (n-1)f''(t) + 2z_0^2 f(t) = 0,$$
(A.1)

$$tg'''(t) - (n-1)g''(t) + 2\bar{z}_0^2 g(t) = 0.$$
(A.2)

Multiplying both sides of (A.1) by g, integrating both sides from 1 to ∞ , and performing integration by parts, we obtain

$$0 = \int_{1}^{\infty} [tf'''g - (n-1)f''g + 2z_{0}^{2}fg]dt$$

$$= tgf''\Big|_{1}^{\infty} - \int_{1}^{\infty} f''(g + g't)dt + \int_{1}^{\infty} [-(n-1)f''g + 2z_{0}^{2}fg]dt$$

$$= \int_{1}^{\infty} [-tf''g' - nf''g + 2z_{0}^{2}fg]dt$$

$$= \int_{1}^{\infty} [-tf''g' + nf'g' + 2z_{0}^{2}fg]dt.$$
(A.3)

Similarly,

$$0 = \int_{1}^{\infty} \left[-tf'g'' + nf'g' + 2\bar{z}_{0}^{2}fg \right] dt.$$
(A.4)

Moreover,

$$\int_{1}^{\infty} [-tf'g'' - tf'g'']dt = -\int_{1}^{\infty} td(f'g')$$

= $-tf'g'\Big|_{1}^{\infty} + \int_{1}^{\infty} f'g'dt$ (A.5)
= $f'(1)g'(1) + \int_{1}^{\infty} f'g'dt.$

Adding (A.3), (A.4) and using (A.5) to simplify the result, we obtain

$$0 = f'(1)g'(1) + (2n+1)\int_{1}^{\infty} f'g'dt + 2(z_0^2 + \bar{z}_0^2)\int_{1}^{\infty} fgdt.$$
 (A.6)

Rearranging (A.6), we have

$$4(y_0^2 - x_0^2) \int_1^\infty |J_n(z_0t)|^2 dt = |z_0|^2 |J_n'(z_0)|^2 + (2n+1)|z_0|^2 \int_1^\infty |J_n'(z_0t)|^2 dt.$$
(A.7)

Since the right side of (A.7) and the integral on its left side are both positive, we must have $y_0^2 - x_0^2 > 0$ and the lemma follows.

Lemma 5. $J_n(z)$ has no zero in $D = \{z \in \overline{\mathbb{C}^+} ||z| > R\}$, where R is sufficiently large.

Proof. Subtracting (A.4) from (A.3), we have

$$0 = \int_{1}^{\infty} t(f'g'' - f''g')dt + 2(z_0^2 - \bar{z}_0^2) \int_{1}^{\infty} fgdt.$$
(A.8)

That is,

$$4x_0y_0\int_1^\infty |J_n(z_0t)|^2 dt = |z_0|^2\int_1^\infty \operatorname{Im}\left(\bar{z}_0tJ_{n-1}(z_0t)J_{n-2}(\bar{z}_0t)\right) dt. \quad (A.9)$$

In the domain D, $J_n(z)$ is well approximated by the leading term of its asymptotic expansion. Let $z_0 = r_0 e^{i\theta_0}$ with $r_0 > 0$ and $\theta_0 \in [-\pi/2, \pi/2]$. Substituting the leading terms of the asymptotic expansions into both sides of (A.9) and simplifying the resulting expressions, we obtain

$$\sin(2\theta_0) \sim -\sin(2\theta_0/3). \tag{A.10}$$

In other words, two sides of (A.9) have opposite sign unless they are both equal to zero, *i.e.*, unless $\theta_0 = 0$ or z_0 is a positive real number. However, $J_n(x) > 0$ when x > 0, as seen from its integral representation (1). And the lemma follows.

Appendix B. The coefficients of Laurent polynomial approximations for J_n

We list the coefficients c_j of Laurent polynomial approximations for evaluating J_n (n = -1, 0, 1, and 2) on each quarter-annulus domain Q_i (i = 1, 2, and 3) in Tables B.4–B.15. That is,

$$J_n(z) \approx \sqrt{\frac{\pi}{3}} \left(\frac{\nu}{3}\right)^{n/2} e^{-\nu} \nu^{N_2} \sum_{j=0}^{N_T - 1} c_j \left(\frac{1}{\nu}\right)^j, \tag{B.1}$$

where $\nu := 3 \left(\frac{z}{2}\right)^{2/3}$. (B.1) is obtained by combining (6) and (21), and rewriting the Laurent polynomial as a power series in $\frac{1}{\nu}$ by pulling out the factor ν^{N_2} .

Table B.4: The coefficients c_j (j = 0, ..., 31) of the Laurent polynomial approximation given by (B.1) to evaluate $J_{-1}(z)$ to 19-digit precision in $Q_1 := \{z \in \mathbb{C} \mid \text{Re}(z) \ge 0, \text{Im}(z) \ge 0, 1 \le |z| \le 3\}$. $N_2 = 10$.

real part	imaginary part
$\begin{array}{c} \mbox{real part} \\ \hline 0.508\ 404\ 632\ 082\ 606\ 781\ 52\ \times\ 10^{-17} \\ -0.745\ 912\ 235\ 026\ 426\ 206\ 60\ \times\ 10^{-14} \\ 0.510\ 342\ 448\ 563\ 248\ 242\ 07\ \times\ 10^{-12} \\ -0.155\ 278\ 534\ 850\ 271\ 007\ 09\ \times\ 10^{-10} \\ 0.264\ 414\ 045\ 122\ 879\ 630\ 95\ \times\ 10^{-9} \\ -0.247\ 487\ 638\ 713\ 530\ 933\ 63\ \times\ 10^{-8} \\ 0.482\ 268\ 582\ 740\ 909\ 041\ 08\ \times\ 10^{-8} \\ 0.216\ 253\ 553\ 725\ 866\ 075\ 08\ \times\ 10^{-6} \\ -0.368\ 717\ 051\ 178\ 481\ 237\ 97\ \times\ 10^{-5} \\ 0.346\ 284\ 048\ 895\ 070\ 301\ 60\ \times\ 10^{-4} \\ 0.999\ 777\ 374\ 590\ 696\ 606\ 94 \\ -0.823\ 278\ 581\ 628\ 169\ 930\ 45\ \times\ 10^{-1} \\ -0\ 745\ 765\ 60\ 41\ 577\ 60\ 41\ 577\ 60\ 41\ 577\ 60\ 41\ 577\ 60\ 41\ 577\ 60\ 41\ 577\ 60\ 41\ 577\ 60\ 41\ 577\ 60\ 41\ 41\ 41\ 577\ 60\ 41\ 577\ 60\ 41\ 41\ 41\ 41\ 41\ 41\ 41\ 41\ 41\ 41$	$\begin{array}{r} {\rm imaginary \ part} \\ \hline -0.174\ 608\ 152\ 994\ 637\ 499\ 48\times 10^{-15} \\ 0.124\ 626\ 002\ 002\ 964\ 530\ 12\times 10^{-13} \\ -0.294\ 298\ 471\ 469\ 682\ 176\ 69\times 10^{-12} \\ 0.263\ 158\ 514\ 306\ 763\ 567\ 96\times 10^{-12} \\ 0.139\ 834\ 751\ 397\ 682\ 449\ 07\times 10^{-9} \\ -0.374\ 213\ 198\ 230\ 171\ 159\ 33\times 10^{-8} \\ 0.543\ 401\ 289\ 326\ 601\ 410\ 72\times 10^{-7} \\ -0.508\ 720\ 998\ 708\ 511\ 613\ 98\times 10^{-6} \\ 0.301\ 340\ 166\ 557\ 595\ 939\ 20\times 10^{-5} \\ -0.739\ 108\ 230\ 704\ 056\ 172\ 19\times 10^{-5} \\ -0.576\ 030\ 836\ 245\ 300\ 251\ 51\times 10^{-4} \\ 0.849\ 768\ 580\ 372\ 614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 372\ 614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 037\ 2614\ 021\ 53\times 10^{-3} \\ 0.694\ 768\ 580\ 572\ 590\ 590\ 50\ 10^{-3} \\ 0.694\ 768\ 580\ 572\ 590\ 590\ 51\ 51\ 50\ 50\ 50\ 50\ 50\ 50\ 50\ 50\ 50\ 50$
$\begin{array}{c} 0.61974789354573766566\times10^{-3}\\ 0.56615182294768079637\times10^{-1}\\ -0.13513677999109029679\\ 0.20971815296188580167\\ -0.13559302399958735143\\ -0.39989888854107271642\\ 0.17398271638333840850\times10^{1}\\ -0.38218064277297175142\times10^{1}\\ -0.38218064277297175142\times10^{1}\\ 0.57404644363343330931\times10^{1}\\ -0.60218557453822568030\times10^{1}\\ 0.38664098293378463250\times10^{1}\\ 0.38664098293378463250\times10^{1}\\ \end{array}$	$\begin{array}{c} -0.60513938190315115982\times10^{-2}\\ 0.30290788934912727755\times10^{-1}\\ -0.11595801511190682178\\ 0.34917394460827128828\\ -0.83031652814884108813\\ 0.15322617865070516148\times10^{1}\\ -0.20720742240832378654\times10^{1}\\ 0.16655645078067718066\times10^{1}\\ 0.36441373700438752895\\ -0.36699889432071554424\times10^{1}\\ 0.65632365306504714202\times10^{1}\\ 0.65632365306504714202\times10^{1}\\ \end{array}$
$\begin{array}{c} -0.25697949139671290942\\ -0.26138368668366285752\times10^1\\ 0.33128062049048194583\times10^1\\ -0.22947550138820496670\times10^1\\ 0.97998841946268481518\\ -0.23273273868918701241\\ 0.13573816724649184659\times10^{-1}\\ 0.64717665310787895482\times10^{-2}\\ -0.11353082240496407813\times10^{-2} \end{array}$	$\begin{array}{l} -0.71730334400727831101\times10^{1}\\ 0.52101729871789289259\times10^{1}\\ -0.22744720239035355439\times10^{1}\\ 0.21824868540130264477\\ 0.43081914671714156557\\ -0.30881694364539401656\\ 0.10103591470624091510\\ -0.16204976273553372187\times10^{-1}\\ 0.91111645509511076869\times10^{-3} \end{array}$

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Table B.5: Similar to Table B.4, c_j (j = 0, ..., 29) for $J_{-1}(z)$ in $Q_2 = \{z \in \mathbb{C} \mid \text{Re}(z) \ge 0, \text{Im}(z) \ge 0, 3 \le |z| \le 15\}$. $N_2 = 0$.

real part	imaginary part
0.99999999999996165301	$0.14180683234758492536\times10^{-12}$
$-0.83333333315888343156 imes 10^{-1}$	$-0.18355475502542401539 imes10^{-10}$
$0.34722202099214306218 imes10^{-2}$	$0.66429512090231781628 imes10^{-9}$
$0.55459217936935525195 imes10^{-1}$	$0.19209070325965982796 imes10^{-7}$
-0.17477009309488548835	$-0.27235872415655243493 imes10^{-5}$
0.47557985079285319878	$0.12329339800149018587 imes10^{-3}$
$-0.12044719601488244381\times10^{1}$	$-0.33379989131254858384 imes10^{-2}$
$0.24160534977076998585 imes10^1$	$0.59920404647268786033 imes10^{-1}$
0.71401934124020221324	-0.69764699651584141417
$-0.60367540682210374145 imes10^2$	$0.38607796239007672158 imes10^1$
$0.60545135048209986187 imes10^3$	$0.32429279475615845398 imes10^2$
$-0.45946367108344566727 imes10^4$	$-0.10841949093356820460 imes10^4$
$0.28358573752155457724 imes10^5$	$0.14766714227455119633 imes10^5$
$-0.13554840952273842275 imes10^6$	$-0.13629699320708838668 imes10^6$
$0.42722335885416276983 imes10^6$	$0.93640538562551055857 imes10^6$
$-0.18717734419017137932 imes10^6$	$-0.49014636853552340206 imes10^7$
$-0.76674698133130508647 imes10^7$	$0.19293996919633486842 imes10^8$
$0.56621620119877490002 imes10^8$	$-0.53379687761932413868 imes10^8$
$-0.24544626054098569413 imes10^9$	$0.76969984306318530811 imes10^8$
$0.73652406083022339655 imes10^9$	$0.11898870845017090857 imes10^9$
$-0.15200293963699011585 imes10^{10}$	$-0.11167744707665950837 imes10^{10}$
$0.18157636339201652460 imes10^{10}$	$0.37001812450286450398 imes10^{10}$
$0.23697005105214074056 imes10^9$	$-0.76973235212132828329 imes10^{10}$
$-0.60541865274209691412 imes10^{10}$	$0.10509437050190981895 imes10^{11}$
$0.13447347591183417529 imes10^{11}$	$-0.82869660102657042317 imes10^{10}$
$-0.16538600326905832899 imes10^{11}$	$0.87889333133786055548 imes10^9$
$0.12108038654949012813 imes10^{11}$	$0.58320160548830925371 imes10^{10}$
$-0.45881323770532016082 imes10^{10}$	$-0.64591210758282531847 imes10^{10}$
$0.33559769561348792357 imes10^9$	$0.30012974946895292083\times10^{10}$
$0.21590442067376607526 imes10^9$	$-0.51553627638896435829\times10^9$

Table B.6: Similar to Table B.4, c_j (j = 0, ..., 19) for $J_{-1}(z)$ in $Q_3 = \{z \in \mathbb{C} \mid \text{Re}(z) \ge 0, \text{Im}(z) \ge 0, 15 \le |z| \le 120\}$. $N_2 = 0$.

real part	imaginary part
$0.10000000000000000211\times10^{1}$	$0.17867305969317471010\times10^{-16}$
$-0.83333333333337062447 imes10^{-1}$	$-0.97723166437483860903 imes 10^{-14}$
$0.34722222219662307873 imes10^{-2}$	$0.18575099531415563550 imes10^{-11}$
$0.55459105063779291569 imes10^{-1}$	$-0.17036760654072501033 imes10^{-9}$
-0.17476652435372606609	$0.72097356819624208386 imes10^{-8}$
0.47552180369947961103	$0.45886722797104214745 imes10^{-7}$
$-0.12045748284986630605 imes10^1$	$-0.23432576523368445886 imes10^{-4}$
$0.24476460069464141708 imes10^1$	$0.14488404302693181855 imes10^{-2}$
-0.19443570247379529707	$-0.51284106279947756551 imes10^{-1}$
$-0.44775070512394599808 imes10^2$	$0.11973398083848165562 imes10^1$
$0.42163459709409223079 imes10^3$	$-0.18732584217083190217 imes10^2$
$-0.30990846226113832847\times10^4$	$0.18064501304516396811 imes10^3$
$0.20913884916390585368 imes10^5$	$-0.59296211153624558148 imes10^3$
$-0.12895996355054874785 imes10^6$	$-0.10821482660282026169\times10^5$
$0.67085295525681611041 imes10^6$	$0.19808862750865237678 imes10^6$
$-0.26337877182586616150 imes10^7$	$-0.16813332113288193078 imes10^7$
$0.67096341894817561042 imes10^7$	$0.85931950401381414532\times10^{7}$
$-0.76571281908120841443 imes10^7$	$-0.26572627858717182234 imes10^8$
$-0.59803448026875748509 imes10^7$	$0.44801684284187004703\times10^8$
$0.19209322347765871037 imes10^8$	$-0.30066013610259277074 imes10^8$

Table B.7: The coefficients c_j (j = 0, ..., 29) of the Laurent polynomial approximation given by (B.1) to evaluate $J_0(z)$ to 19-digit precision in $Q_1 := \{z \in \mathbb{C} \mid \text{Re}(z) \ge 0, \text{Im}(z) \ge 0, 1 \le |z| \le 3\}$. $N_2 = 11$.

real part	imaginary part
$-0.90832607641433626723\times10^{-16}$	$-0.12971716857438253177 \times 10^{-15}$
$0.12389804620230878343 \times 10^{-12}$	$0.12374086345769475560 \times 10^{-12}$
$0.18925665936446973863\times10^{-11}$	$-0.43594042425001909808 \times 10^{-11}$
$-0.93436124699082728782 \times 10^{-11}$	$0.71610240171455947215\times10^{-11}$
$0.21005471373426356192\times10^{-5}$	$-0.31570720100508497322 \times 10^{-10}$
$-0.27921469604412283831 \times 10^{-8}$	$-0.10267754766776108185 \times 10^{-8}$
$0.22210697286892781643 imes10^{-7}$	$0.25281961277933484578 imes10^{-7}$
$-0.71494613675879873372 \times 10^{-7}$	$-0.30771058781715872427 \times 10^{-6}$
$-0.63954539217286436591 imes10^{-6}$	$0.24264961215641857661 imes10^{-5}$
$0.11477599934330755236 imes10^{-4}$	$-0.12692783111400141826 \times 10^{-4}$
$-0.94447601385518738118 imes10^{-4}$	$0.36582122360442313063 \times 10^{-4}$
$0.10005227131416389419 imes10^1$	$0.42527679632933790645 \times 10^{-4}$
$-0.85416339695541973974 imes10^{-1}$	$-0.11668274275597700974 imes10^{-2}$
$0.92680088758957499003 imes10^{-1}$	$0.75525130431164108091 imes10^{-2}$
-0.12830962183399732914	$-0.31999279672322587569 imes10^{-1}$
0.18262801902460105636	0.10112601789451786253
-0.22086524963618477505	-0.24794786230219646677
0.16700979504519707012	0.47574461283438653640
$0.62230841342939847177 imes10^{-1}$	-0.70546898276037268906
-0.46160590435660301333	0.77405336431438665012
0.86282905145667261861	-0.54883285003677036633
$-0.10157725409050061178 imes10^1$	$0.89653554833115686588 imes10^{-1}$
0.81127042933258073193	0.34223426976661392785
-0.40588870780744941328	-0.50414950527612358637
$0.70646295770216518095 imes10^{-1}$	0.39020053902964327900
$0.62576278207207491689 imes10^{-1}$	-0.18665248819340615755
$-0.56069463746156990540 imes10^{-1}$	$0.51369345024189480453 imes10^{-1}$
$0.20854565059593331201 imes10^{-1}$	$-0.49622460298807761080 \times 10^{-2}$
$-0.37736411371630856848 imes10^{-2}$	$-0.10706620427594972634 imes10^{-2}$
$0.24658062816300990462\times 10^{-3}$	$0.24793548656986345025\times10^{-3}$

Table B.8: Similar to Table B.7, c_j (j = 0, ..., 29) for $J_0(z)$ in $Q_2 := \{z \in \mathbb{C} \mid \text{Re}(z) \ge 0, \text{Im}(z) \ge 0, 3 \le |z| \le 15\}$. $N_2 = 0$.

real part	imaginary part
0.99999999999988637217	$-0.86635400939375232846 \times 10^{-13}$
$-0.83333333323131499972 imes 10^{-1}$	$0.22519321671024025250 imes10^{-10}$
$0.86805555689429826898 imes10^{-1}$	$-0.20734497580441697938 imes10^{-8}$
-0.11815206514175737947	$0.96339270441630372077 imes10^{-7}$
0.17969333057187777654	$-0.23057043651727157918 imes10^{-5}$
-0.24337342169790375842	$0.98098354524518453957 imes10^{-5}$
$0.14764429473872620763 imes10^{-1}$	$0.12782321450918857049 imes10^{-2}$
$0.23309627937902507555 imes10^1$	$-0.51073707172260220779 \times 10^{-1}$
$-0.16480981490288764106 imes10^2$	$0.11291585827876199238 imes10^1$
$0.92305117544447684520 imes10^2$	$-0.17262096251706180600 imes10^2$
$-0.51600417922753718015 imes10^3$	$0.19328752764190907274 imes10^3$
$0.32464931691139105961 imes10^4$	$-0.15901053336317843819 imes10^4$
$-0.22314060512920598345 imes10^5$	$0.90599663123843335441 imes10^4$
$0.14604708231321899614 imes10^6$	$-0.26233323039837437629 imes10^{5}$
$-0.81145513740976803311 imes10^{6}$	$-0.98104807756320344084 imes10^{5}$
$0.35457971304583219348 imes10^7$	$0.18294948704040264718 imes10^7$
$-0.11084057351725538629 imes10^8$	$-0.13170160809055557556 imes10^8$
$0.17946742587727355394 imes10^8$	$0.62986605485232028006 imes10^8$
$0.35800094666248698896 imes10^8$	$-0.21614451081432945159 imes10^9$
$-0.37304955427569800721 imes10^9$	$0.52194699072674118229 imes10^9$
$0.14479420488451058753 imes10^{10}$	$-0.75867692917492941354 imes10^9$
$-0.35951788816564166863 \times 10^{10}$	$0.22618398696580184017 imes10^8$
$0.60001115021506662033 imes10^{10}$	$0.31006219604826621451 imes10^{10}$
$-0.60780548332669974319 imes10^{10}$	$-0.87891548726761143921 \times 10^{10}$
$0.15576242961336566425 imes10^{10}$	$0.13885136549901956758 imes10^{11}$
$0.55250083598538675266 imes10^{10}$	$-0.13622496533575828228\times10^{11}$
$-0.92625760934489411655 \times 10^{10}$	$0.75732780185372212772 \times 10^{10}$
$0.69543035636142439747 imes10^{10}$	$-0.12850366647773799533 imes 10^{10}$
$-0.25630570764916110006 imes10^{10}$	$-0.85541751450691424086 \times 10^{9}$
$0.33863352297553708594 imes10^9$	$0.36939690751181780209 imes10^9$

Table B.9: Similar to Table B.7, c_j (j = 0, ..., 19) for $J_0(z)$ in $Q_3 := \{z \in \mathbb{C} \mid \text{Re}(z) \ge 0, \text{Im}(z) \ge 0, 15 \le |z| \le 120\}$. $N_2 = 0$.

real part	imaginary part
0.99999999999999996930 = 0.833 333 333 333 199 009 02 × 10 ⁻¹	$0.17593864033911935746\times10^{-16}$ -0.16846147898292977950×10 ⁻¹⁵
$0.86805555553423816181\times10^{-1}$	$-0.11453394704160148063\times10^{-11}$
-0.11815200602347672298	$0.23048385969290284678 \times 10^{-9}$
0.17968950772370040285 - 0.243 237 776 508 143 247 72	$-0.21945385998958748345 \times 10^{-1}$ 0.120.879,546,268,738,763,83 $\times 10^{-5}$
$0.11700570395152938411\times10^{-1}$	$-0.38103854798729283635 \times 10^{-4}$
$0.23745206848674586241\times10^{1}$	$0.41346400173493025458 imes10^{-3}$
$-0.16758837066520389957 imes10^2$	$0.20437630577696960942 \times 10^{-1}$
$0.88966545765421650942\times10^{2}$	$-0.12064317238786081373\times10^{1}$
$-0.39494374005352553520 \times 10$ 0.14059897264393353350 × 10 ⁴	$-0.62474820464709927286 \times 10^{3}$
$-0.40451355322086019935 imes10^4$	$0.80876358650678996996 imes10^4$
$0.19384543104359811933\times10^5$	$-0.74559599790322239386\times 10^5$
$-0.20517688352680136030 imes10^6$	$0.48086289900895412320 imes 10^6$
$0.17043521840830888384 imes10^{\prime}$	$-0.20518647360918521409 \times 10^{7}$
-0.87705077033153779372 imes10'	$0.49903982905075629408\times10^{\prime}$
$0.20797195290998643893 \times 10^{\circ}$	$-0.308(4807345514436382 \times 10^{\circ})$
$-0.43446793431180983346 \times 10^{-1}$ 0.26603074410888988636 $\times 10^{8}$	$-0.14196190720385577590 \times 10^{-1}$ 0.25674511583029811722 $\times 10^{8}$
0.200 000 111100 000 000 00 × 10	0.200 10 110 000 200 111 22 × 10

Table B.10: The coefficients c_j (j = 0, ..., 29) of the Laurent polynomial approximation given by (B.1) for evaluation of $J_1(z)$ to 19-digit precision in $Q_1 := \{z \in \mathbb{C} \mid \text{Re}(z) \ge 0, \text{Im}(z) \ge 0, 1 \le |z| \le 3\}$. $N_2 = 11$.

real part	imaginary part
$0.11005198342846485755 \times 10^{-15}$	$-0.69497479897694901798\times10^{-16}$
$-0.10253717390952836750\times10^{-13}$	$0.57344733061298400754 imes10^{-15}$
$0.35541980746147250213 imes10^{-12}$	$0.17138035947739436987 \times 10^{-12}$
$-0.57065663426773316925 \times 10^{-11}$	$-0.80097752680384861353 imes10^{-11}$
$0.21377942402322032801 imes 10^{-10}$	$0.17745761143488866634 imes10^{-9}$
$0.92504177773563681659 imes10^{-9}$	$-0.23491690282446912066 imes10^{-8}$
$-0.21958132206773784869 \times 10^{-7}$	$0.18715255541162079236 imes10^{-7}$
$0.26744146813435088020 imes10^{-6}$	$-0.60538192727561566571 imes10^{-7}$
$-0.21417964338884026393 imes10^{-5}$	$-0.55166777471060148467 imes10^{-6}$
$0.11605229803403698059 imes10^{-4}$	$0.10089206827157776627 imes10^{-4}$
$-0.36902286245955926331 imes 10^{-4}$	$-0.85734523630464858844 imes10^{-4}$
0.99998886026520537047	$0.49984760395796500059 imes10^{-3}$
0.41764834336748872867	$-0.21759152835287886250 imes10^{-2}$
-0.12880435876254801279	$0.72384623466304053406 \times 10^{-2}$
0.10001797075393494104	$-0.18204681378409207336 imes10^{-1}$
-0.12326994533416519599	$0.32494324768891981532 imes10^{-1}$
0.19390815724013910177	$-0.31048056008057556412 imes10^{-1}$
-0.30446220948476603076	$-0.28094371375086694729 imes10^{-1}$
0.40217178304394830919	0.18701432829068936487
-0.39045411378013014641	-0.42965195739247365992
0.20299486067051564492	0.64270959413495860386
0.10086985423793876517	-0.67712707054748086683
-0.34706311627097173040	0.48943085630911667832
0.39717127772830854281	-0.20388651987798637536
-0.27627970190658614495	$-0.36920639378264876881 \times 10^{-2}$
0.12093911370868755832	$0.67168398413971524804 \times 10^{-1}$
$-0.28845724988884421284\times10^{-1}$	$-0.45442309954198940539\times10^{-1}$
$0.97253793671434469328 imes10^{-3}$	$0.15236745276873557399 imes10^{-1}$
$0.11989911484170176670 imes10^{-2}$	$-0.25400926250086296020 \times 10^{-2}$
$-0.20436565348663071365 imes10^{-3}$	$0.14672057879876671250 imes10^{-3}$

Table B.11: Similar to Table B.10, c_j (j = 0, ..., 29) for $J_1(z)$ in $Q_2 := \{z \in \mathbb{C} \mid \text{Re}(z) \ge 0, \text{Im}(z) \ge 0, 3 \le |z| \le 15\}$. $N_2 = 0$.

real part	imaginary part
$0.10000000000001559822\times10^{1}$	$-0.64432580975613771082\times10^{-13}$
0.41666666663765045792	$-0.30929290380316585308 imes10^{-11}$
-0.12152777575602031881	$0.13847887495450779512 imes10^{-8}$
$0.64139599010730684601 imes10^{-1}$	$-0.11804143890505963878 imes10^{-6}$
$0.19340333876868250914 imes10^{-1}$	$0.52525129103659222882 imes 10^{-5}$
-0.31085396288117458325	$-0.14209942615180352958 imes10^{-3}$
$0.14076112393497740595 imes10^1$	$0.22843270818239797021 imes 10^{-2}$
$-0.53034603582858591137\times10^{1}$	$-0.12508426902104267053 imes10^{-1}$
$0.16843969837042581097 imes10^2$	-0.39633859151173122745
$-0.30896103962008743537 imes10^2$	$0.13366867242921989470 imes10^2$
$-0.12354078658896338072 imes10^3$	$-0.23100296557839310161 imes10^3$
$0.16486028729189802772 imes10^4$	$0.27798011316789725915 imes10^4$
$-0.79365600281962572617 imes10^4$	$-0.25085379438345641690 imes10^5$
$-0.17876338162032969792 imes10^4$	$0.17337711302201102756 imes10^6$
$0.36042195374432098049 imes10^6$	$-0.90932303898462350515 imes10^6$
$-0.33818329308700649502 imes10^7$	$0.34283955443872702745 imes10^7$
$0.19449202897749494834 imes10^8$	$-0.74625499598352743942 imes10^7$
$-0.79047083372251398893 imes10^8$	$-0.61704036445636743642 imes10^7$
$0.22787889419550420217 imes10^9$	$0.13542454969531036697 imes10^9$
$-0.41550638595000604177 imes10^9$	$-0.65483373280300662671 imes10^9$
$0.17290712369617688112 imes10^9$	$0.19664895233379210138 imes10^{10}$
$0.16763777971399520871 imes10^{10}$	$-0.40005075227615193294 imes10^{10}$
$-0.63097117018939634689 imes10^{10}$	$0.51442671230430695393 imes10^{10}$
$0.12639623708614918540 imes10^{11}$	$-0.24250533919233200087 imes10^{10}$
$-0.16021424838185937866 imes10^{11}$	$-0.51063223260105245852 imes10^{10}$
$0.12194470039177973074 imes10^{11}$	$0.12801295658115048467 imes10^{11}$
$-0.36617569740928099053 imes10^{10}$	$-0.13906995324951658335 imes10^{11}$
$-0.20882430849265640704 imes10^{10}$	$0.82351413910371916318 imes10^{10}$
$0.22256324759689985206 imes10^{10}$	$-0.23606999259817906485 imes10^{10}$
$-0.57357275466466452587 imes10^9$	$0.18100167963268264995 imes10^9$

Table B.12: Similar to Table B.10, c_j (j = 0, ..., 19) for $J_1(z)$ in $Q_3 := \{z \in \mathbb{C} \mid \text{Re}(z) \ge 0, \text{Im}(z) \ge 0, 15 \le |z| \le 120\}$. $N_2 = 0$.

real part	imaginary part
$0.10000000000000000088\times 10^{1}$	$-0.37104682094436741073 imes 10^{-16}$
0.41666666666665693268	$0.10633105786560679943 \times 10^{-13}$
-0.12152777777532530135	$-0.81783483309751920964 \times 10^{-12}$
$0.64139660206844395055 imes10^{-1}$	$-0.53206746309599215652\times10^{-10}$
$0.19340376146506349534 imes10^{-1}$	$0.14714309085259108266 \times 10^{-7}$
-0.31092901473600760433	$-0.12847633757844741159 imes10^{-5}$
$0.14108230204421526148 imes10^1$	$0.64029804309267135469 imes10^{-4}$
$-0.53814287035192935174 imes10^1$	$-0.19729163937458920069 imes10^{-2}$
$0.18099040506545928510 imes10^2$	$0.33959266046710551528 imes10^{-1}$
$-0.44222521101771005259 imes10^2$	$-0.46443129149364017364 imes10^{-1}$
$-0.49281712505609892221\times10^2$	$-0.14673338698274539790 imes10^2$
$0.19120731039799297866 imes10^4$	$0.44654655766491527724 imes10^3$
$-0.18480296407139556113 imes10^5$	$-0.76529257176182418914 imes10^4$
$0.11949451283301751133 imes10^6$	$0.88037548115504996527 imes10^5$
$-0.51411364057002663824 imes10^6$	$-0.70563400281206830770 imes10^6$
$0.10973535392819828712 imes10^7$	$0.39099452669156576703 imes10^7$
$0.17331435565199135031 imes10^7$	$-0.14355007661133573735 imes10^8$
$-0.19127195748032646177 imes10^8$	$0.31757253245227946371\times10^8$
$0.50801574443329963657 imes10^8$	$-0.33700657946301331333 imes10^8$
$-0.47910288059234253994 imes10^8$	$0.61171907037609011958 imes10^7$

Table B.13: The coefficients c_j (j = 0, ..., 31) of the Laurent polynomial approximation given by (B.1) to evaluate $J_2(z)$ to 19-digit precision in $Q_1 = \{z \in \mathbb{C} \mid \operatorname{Re}(z) \ge 0, \operatorname{Im}(z) \ge 0, 1 \le |z| \le 3\}$. $N_2 = 11$.

real part	imaginary part
$0.31866632685819612221 \times 10^{-16}$	$0.62278738969830137987 \times 10^{-16}$
$0.23374202488114714431 \times 10^{-10}$	$-0.58368186829092031391 \times 10^{-11}$
$-0.12464763262921601144 \times 10^{-12}$	$0.20259199528722770999 \times 10^{-12}$
$0.54938914867389395195 \times 10^{-11}$	$-0.30688169770355401691 \times 10^{-11}$
$-0.12158575842106281373 imes10^{-9}$	$0.20706119567919592298 imes10^{-12}$
$0.16018003273928560165 imes10^{-8}$	$0.88216842473436341238 imes10^{-9}$
$-0.11941475881449767865 imes10^{-7}$	$-0.18815539800292924303 imes10^{-7}$
$0.15306868790345339446 imes10^{-7}$	$0.22567034551573649710 imes10^{-6}$
$0.80407254481241384544 imes10^{-6}$	$-0.17772860498030423720 imes10^{-5}$
$-0.11465040286971344849 imes10^{-4}$	$0.88993097363922792729 imes10^{-5}$
$0.92698417450695281624 imes10^{-4}$	$-0.17153044258356809874 imes10^{-4}$
0.99948171991549382451	$-0.15046194458696453245 imes10^{-3}$
$0.14187062252412931802 imes10^1$	$0.18349141691729387994 imes10^{-2}$
-0.12647835873372535821	$-0.11428276705995908659 imes10^{-1}$
0.18603598111474663325	$0.50646899433597524505 imes10^{-1}$
-0.28764948748580855321	-0.17268243476553171463
0.37671933372380252436	0.46496076480920682474
-0.28673562730306685162	-0.99215999495574618374
-0.25260802503698437799	$0.16511988403580568251 imes10^{1}$
$0.14171261472971538418 imes10^1$	$-0.20381244501795426809 imes10^1$
$-0.29457967249708858073 imes10^1$	$0.15782077203635430607 imes10^1$
$0.40284406518941590462 imes10^1$	$-0.37723354335564292819 imes10^{-1}$
$-0.38128950013920692643 imes10^1$	$-0.19907064876530093167 imes10^1$
$0.22475249180098600179 imes10^1$	$0.33343987769848875142 imes10^1$
-0.29845507995469994980	$-0.32351706541913757406 imes10^1$
-0.89544287444481992984	$0.20454467870089365942 imes10^1$
$0.10175639567579240708 imes10^1$	-0.77474351850574639868
-0.58788778865800093021	$0.83503119281933322786 \times 10^{-1}$
0.20009319020081182763	$0.77689730095477593636 imes10^{-1}$
$-0.35965432090839096074 imes10^{-1}$	$-0.43797603367743883119 imes10^{-1}$
$0.16769420558534530117 imes10^{-2}$	$0.95756365430182522746\times10^{-2}$
$0.27256710553195448121 imes10^{-3}$	$-0.76612682266254092889 imes10^{-3}$

Table B.14: Similar to Table B.13, c_j (j = 0, ..., 29) for $J_2(z)$ in $Q_2 = \{z \in \mathbb{C} \mid \text{Re}(z) \ge 0, \text{Im}(z) \ge 0, 3 \le |z| \le 15\}$. $N_2 = 0$.

real part	imaginary part
0.99999999999998061808	$0.16248252113309315678\times10^{-12}$
$0.14166666666829275403 imes10^1$	$-0.23049278428420443693 imes 10^{-10}$
-0.12152777988809406380	$0.10578621411215754890 imes10^{-8}$
0.18566756593083182116	$0.28424278401127180014 imes10^{-8}$
-0.35199891174698963256	$-0.24190974576544113333 imes 10^{-5}$
0.74514028477189295514	$0.12627505833721768389 imes10^{-3}$
$-0.15698688665407300795 imes10^1$	$-0.36906906699773089097 \times 10^{-2}$
$0.24402663496436078603 imes10^1$	$0.71150880132384387316 imes10^{-1}$
$0.42564778384044014501 imes10^1$	-0.92028028234742281569
$-0.86628781485086884613 imes10^2$	$0.69659670041612025801 imes10^1$
$0.76780431897157477513 imes10^3$	$0.12336022855791130638 imes10^1$
$-0.56054291413858448509 imes10^4$	$-0.86886243977848514246 imes10^3$
$0.34709760011543563256 imes10^{5}$	$0.13997432791721954713 imes10^{5}$
$-0.17268737514629848298 imes10^6$	$-0.13922909431496294516 imes10^6$
$0.61187840164952315907 imes10^6$	$0.10063302790914770506 imes10^7$
$-0.90229280253909143653 imes10^6$	$-0.55090070720520207541 imes10^7$
$-0.58239693510024439315 imes10^7$	$0.22810067701144771236 imes10^8$
$0.55585305303478347757 imes10^8$	$-0.68226520995783255732 imes10^8$
$-0.26301929032497527308 imes10^9$	$0.12293641811239037282 imes10^9$
$0.84052909616481142218 imes10^9$	$0.21196827324494884533 imes10^8$
$-0.18657792154606295188 \times 10^{10}$	$-0.10131267803607086455 imes10^{10}$
$0.25915388904113076082 \times 10^{10}$	$0.38427156992160196076 imes10^{10}$
$-0.92517922470966062167\times10^{9}$	$-0.86034302033451782637\times10^{10}$
$-0.51036926935170978903 \times 10^{10}$	$0.12626524314344868926\times10^{11}$
$0.13657170819172273067 \times 10^{11}$	$-0.11301642152404259095\times10^{11}$
$-0.18234815409451008101 imes10^{11}$	$0.35546723514474331442 \times 10^{10}$
$0.14355121657220596589 \times 10^{11}$	$0.45839423517954065893 \times 10^{10}$
$-0.61032052837541757328\times10^{10}$	$-0.64509344322401839565\times10^{10}$
$0.84275175499628369228 \times 10^{9}$	$0.32773844995198787342 \times 10^{10}$
$0.15857079566801991882 imes 10^{3}$	$-0.60570352545416858098\times10^{3}$

Table B.15: Similar to Table B.13, c_j (j = 0, ..., 19) for $J_2(z)$ in $Q_3 = \{z \in \mathbb{C} \mid \text{Re}(z) \ge 0, \text{Im}(z) \ge 0, 15 \le |z| \le 120\}$. $N_2 = 0$.

real part	imaginary part
$0.10000000000000000268 \times 10^{1}$ 0 1/1 666 666 666 666 076 32 × 10^{1}	$0.16274386801955295004 \times 10^{-16}$ -0.10287883462694876761 × 10 ⁻¹³
-0.12152777777773596334	$0.21278064340420781315\times10^{-11}$
0.18566743838222558548 -0 351 994 534 212 129 746 48	$-0.21346739588855205915 imes10^{-9}$ $010810020949854670351 imes10^{-7}$
0.745 056 201 785 816 040 77	$-0.12758335415838672536\times10^{-6}$
$\begin{array}{c} -0.15694400222283724037\times10^{1} \\ 0.24655058999322356159\times10^{1} \end{array}$	$-0.19194750913862712585\times10^{-4}$ $0.14693073094005325564\times10^{-2}$
$0.33607089142515892731\times10^1$	$-0.57088317524506061035 imes10^{-1}$
$-0.69853332457831300099\times10^{-}\\0.55610818499114251623\times10^{3}$	$\begin{array}{c} 0.14378014381758075690\times10^{2} \\ -0.24643410535733376707\times10^{2} \end{array}$
$-0.37382459971986918636 \times 10^4$	$0.27993273451335282871 imes10^3$ $0.17720520252252827568 imes10^4$
$-0.14415275031702837209\times10^{6}$	$-0.95697379334253398976\times10^3$
$0.75853591811873756981 imes10^{6}$ 0.200.728.616.687.601.455.28 \times 10 ⁷	$0.14271567167383035412 \times 10^{6}$ 0.150.020.000.512.270.200.84 × 10 ⁷
$-0.30972801008709145538 \times 10$ 0.85376658237794416873 × 10 ⁷	$-0.15005099051227929984\times10$ $0.84572529844164518814\times10^7$
$-0.12288668565761833934\times10^{8}$	$-0.27945310142205107664\times10^{8}$ 0,400,727,827,887,756,884,47 × 10 ⁸
$0.16376853579174702084\times10^{8}$	$-0.35949830322064479962\times10^{8}$

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