Numerical Validation of a Boundary Element Method With $\boldsymbol{E}$ and $\frac{\partial \boldsymbol{E}}{\partial N}$ as the Boundary Unknowns

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***Abstract*—We recently developed a surface integral equation method where the electric field and its normal derivative are chosen as the boundary unknowns. After reviewing this formulation, we present preliminary numerical calculations that show good agreement with the known results. These calculations are encouraging and invite the further development of the numerical solution.**

I. Introduction

We have recently formulated a frequency domain surface integral equation method [1] that is applicable to penetrable closed surface scatterers. The method has several unique applications and advantages over the standard Stratton–Chu formulation as discussed in [1]. In our formulation, we choose the electric field (E-field) and its normal derivative as the boundary unknowns. This choice leads to 12 scalar unknowns on the surface of the scatterer; for each homogeneous region we have three scalar unknowns associated with the E-field and three scalar unknowns associated with its normal derivative. Similar to a typical surface integral equation formulation, our formulation is also based on the Green's theorem (Green's second identity). This formulation leads to *six* scalar equations, and thus it must be supplemented with *six* additional constraints in order to have the same number of equations as unknowns. Three of these constraints come from the well-known continuity condition of the E-field across an interface and the other three come from the recently derived continuity condition for the normal derivative of the E-field [1]–[3].

In this paper, we numerically solve the above discussed equations for spherical scatterers and compare the results to the well-known Mie series solution. We also comment on the choice of the basis functions in the Galerkin's method and its effects on numerical convergence.

II. Formulation Review

Consider a scatterer with permittivity $\overset{2}{\epsilon}$ and permeability $\overset{2}{\mu}$. The space surrounding the scatterer is assumed to be lossless with permittivity $\overset{1}{\epsilon}$ and permeability $\overset{1}{\mu}$, i.e., $\{\overset{1}{\epsilon}, \overset{1}{\mu}\} \in \mathbb{R}$. If we apply the Green's second identity to the scatterer and the surrounding space, then, after setting the observation point on the surface of the scatterer, we obtain:

$$\overset{\text{inc}}{\boldsymbol{E}}(\widetilde{S}) - \fint_\Sigma \left[\overset{1}{G}\frac{\partial \overset{1}{\boldsymbol{E}}}{\partial N} - \overset{1}{\boldsymbol{E}}\frac{\partial \overset{1}{G}}{\partial N}\right] \mathrm{d}S = \frac{1}{2}\overset{1}{\boldsymbol{E}}(\widetilde{S}), \tag{1a}$$

and

$$\fint_\Sigma \left[\overset{2}{G}\frac{\partial \overset{2}{\boldsymbol{E}}}{\partial N} - \overset{2}{\boldsymbol{E}}\frac{\partial \overset{2}{G}}{\partial N}\right] \mathrm{d}S = \frac{1}{2}\overset{2}{\boldsymbol{E}}(\widetilde{S}), \tag{1b}$$

where $\overset{\text{inc}}{\boldsymbol{E}}$ is the incident E-field, $\fint$ denotes the Cauchy principal value integral, Σ denotes the surface of the scatterer, $\frac{\partial}{\partial N}$ denotes the normal derivative, G is the free-space Green's function, and $\widetilde{S}$ is the observation point on Σ. In (1), the overset digit indicates if the quantity is associated with the scatterer or the surrounding space, e.g., $\overset{2}{\boldsymbol{E}}$ is the E-field just inside the scatterer. In the Gaussian unit system, the continuity condition for the E-field across an interface can be written as [1]:

$$\overset{2}{\boldsymbol{E}} = \overset{\text{r}}{\epsilon}^{-1}\left(\boldsymbol{N}\cdot\overset{1}{\boldsymbol{E}}\right)\boldsymbol{N} + \left(\boldsymbol{S}^\alpha\cdot\overset{1}{\boldsymbol{E}}\right)\boldsymbol{S}_\alpha, \tag{2a}$$

and the continuity condition for its normal derivative as [1]:

$$\begin{aligned}\frac{\partial \overset{2}{\boldsymbol{E}}}{\partial N} = \overset{\text{r}}{\mu}\left(\frac{\partial \overset{1}{\boldsymbol{E}}}{\partial N} - \nabla^\alpha\left[\left(\boldsymbol{N}\cdot\overset{1}{\boldsymbol{E}}\right)\boldsymbol{S}_\alpha - \left(\boldsymbol{S}_\alpha\cdot\overset{1}{\boldsymbol{E}}\right)\boldsymbol{N}\right]\right)\\ + \nabla^\alpha\left[\left(\boldsymbol{N}\cdot\overset{2}{\boldsymbol{E}}\right)\boldsymbol{S}_\alpha - \left(\boldsymbol{S}_\alpha\cdot\overset{2}{\boldsymbol{E}}\right)\boldsymbol{N}\right],\end{aligned} \tag{2b}$$

where $\overset{\text{r}}{\mu} = \overset{2}{\mu}/\overset{1}{\mu}$, $\overset{\text{r}}{\epsilon} = \overset{2}{\epsilon}/\overset{1}{\epsilon}$, $\boldsymbol{N}$ is the unit-normal pointing out of the scatterer, $\boldsymbol{S}_\alpha$ is the surface covariant basis [4], and ∇^α is the contravariant surface derivative [4]. Notice that (2) is written in the Einstein summation convention where the Greek indices range from 1 to 2. Substituting (2) into (1b) and using Gauss's theorem in two dimensions yields [1]:

$$\begin{aligned}\frac{1}{2}\overset{2}{\boldsymbol{E}}(\widetilde{S}) = &\fint_\Sigma\left[\overset{\text{r}}{\mu}\overset{2}{G}\frac{\partial \overset{1}{\boldsymbol{E}}}{\partial N} - \overset{1}{\boldsymbol{E}}\frac{\partial \overset{2}{G}}{\partial N}\right]\mathrm{d}S\\ &+ \left(\overset{\text{r}}{\mu} - \overset{\text{r}}{\epsilon}^{-1}\right)\fint_\Sigma\left(\boldsymbol{N}\cdot\overset{1}{\boldsymbol{E}}\right)\boldsymbol{\nabla}\overset{2}{G}\,\mathrm{d}S\\ &+ \left(1 - \overset{\text{r}}{\mu}\right)\fint_\Sigma\left(\overset{1}{\boldsymbol{E}}\cdot\boldsymbol{\nabla}\overset{2}{G}\right)\boldsymbol{N}\,\mathrm{d}S,\end{aligned} \tag{3a}$$

where

$$\overset{2}{\boldsymbol{E}} = \overset{1}{\boldsymbol{E}} + \left(\overset{\text{r}}{\epsilon}^{-1} - 1\right)\left(\boldsymbol{N}\cdot\overset{1}{\boldsymbol{E}}\right)\boldsymbol{N}. \tag{3b}$$

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Equation (3) and (1a) form a set of *six* scalar integral equations with *six* scalar unknowns, namely, $\overset{1}{\boldsymbol{E}}$ and $\frac{\partial}{\partial N}\overset{1}{\boldsymbol{E}}$. This is the set of the integral equations that we numerically solve in the next section.

III. Numerical Calculations

We discretize the spherical scatterer with flat triangular elements and construct a basis for the E-field and its normal derivative. We use piecewise constant basis functions for each component associated with the triangle surfaces. Thus, the number of unknowns is six times the number of the triangular elements. Furthermore, we use Galerkin's method to discretize the equations. In other words, the test and basis functions are identical. It is worth noting that the basis functions *do not* enforce any continuity conditions for the E-field or its normal derivative along the surface. Hence, it is clear that we cannot obtain an optimal convergence rate. Moreover, we anticipate that the sharp wedges may also cause some difficulties. Finding a better set of basis functions is an interesting question for future research.

The integral equation set given by (3) and (1a) contains strongly singular integrals. The gradient of the Green's function has the strongest singularity and we decompose it into the normal and surface derivative parts. With the help of integration by parts, the latter one reduces to an integral over a triangle surface and a closed integral over the triangle's edges. We evaluate these integrals using the standard singularity extraction technique [5] in which the singular part is calculated analytically and the remaining part is calculated numerically.

To assess the method, we compare the radar cross section (RCS) of a sphere in free-space meshed by 940 flat triangular patches with the Mie series solution. Fig. 1 shows the RCS of a dielectric sphere with $\overset{1}{k}\rho = 1$, $\overset{2}{\epsilon} = 4$, and $\overset{2}{\mu} = 1$, where ρ is the radius of the sphere and Fig. 2 shows the RCS of a lossy sphere with $\overset{1}{k}\rho = 4$, $\overset{2}{\epsilon} = -2 + \mathrm{i}$, and $\overset{2}{\mu} = 1$. From the figures, we see that our solution agrees well with the Mie series solution in both the dielectric case and the lossy case. More specifically, the L^2-norm relative error of the far-field $\|\boldsymbol{E}\|^2$ integrated over a solid angle is 4.832×10^{-3} for Fig. 1 and 9.360×10^{-3} for Fig. 2.

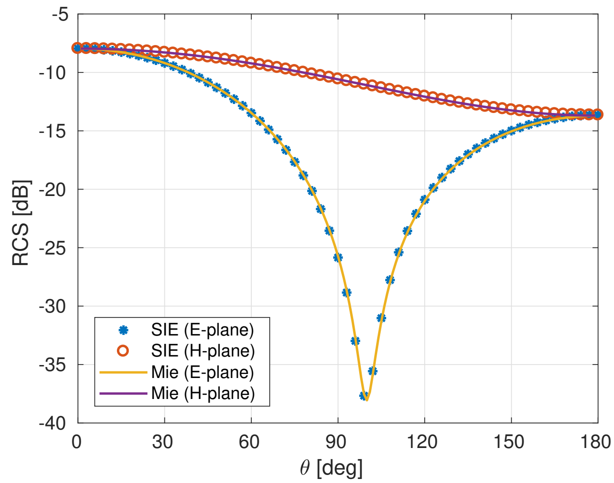

Fig. 1. (Color online) Comparison of the dielectric sphere's RCS as a function of the scattering angle θ computed via the surface integral equation (SIE) method with the Mie series solution.

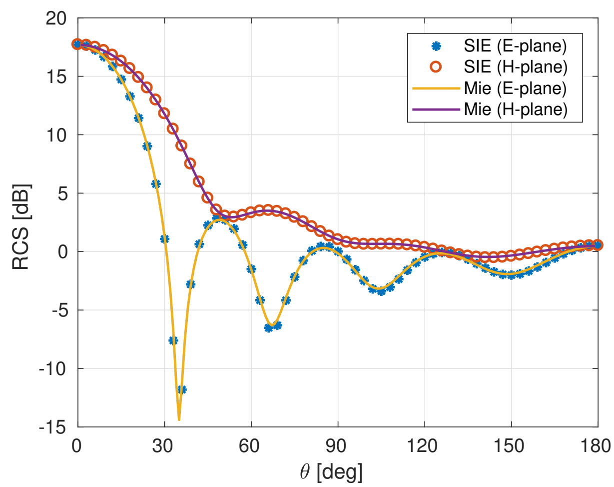

Fig. 2. (Color online) Comparison of the lossy sphere's RCS as a function of the scattering angle θ computed via the surface integral equation (SIE) method with the Mie series solution.

IV. Conclusions

We numerically tested a recently formulated surface integral equation method where the electric field and its normal derivative are chosen as the boundary unknowns. The preliminary results presented here are in agreement with the Mie series solution for both dielectric and lossy spheres. Furthermore, the method seems to be viable for numerical computations and may be further improved if we employ basis functions that enforce the continuity conditions.

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