

Evolution of Ultrafine Particle Size Distributions Following Indoor Episodic Releases: Relative Importance of Coagulation, Deposition and Ventilation

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ABSTRACT

Indoor ultrafine particles (UFP, < 100 nm) undergo aerosol processes such as coagulation and deposition, which alter UFP size distribution and accordingly the level of exposure to UFP of different sizes. This study investigates the decay of indoor UFP originated from five different sources: a gas stove and an electric stove, a candle, a hair dryer, and power tools in a residential test building. An indoor aerosol model was developed to investigate differential effects of coagulation, deposition and ventilation. The coagulation model includes Brownian, van der Waals and viscosity forces, and also fractal geometry for particles > 24 nm. The model was parameterized using different values of the Hamaker constant for predicting the coagulation rate. Deposition was determined for two different conditions: central fan on vs. central fan off. For the case of a central fan running, deposition rates were measured by using a nonlinear solution to the mass balance equation for the whole building. For the central fan off case, an empirical model was used to estimate deposition rates. Ventilation was measured continuously using an automated tracer gas injection and sampling system. The study results show that coagulation is a significant aerosol process for UFP dynamics and the primary cause for the shift of particle size distribution following an episodic high-concentration UFP release with no fans on.. However, with the central mechanical fan on, UFP deposition loss is substantial and comparable to the coagulation loss. These results suggest that coagulation should be considered during high concentration periods ($> 20,000 \text{ cm}^{-3}$), while particle deposition should be treated as a major loss mechanism when air recirculates through ductwork or mechanical systems.

Keywords: Ultrafine particles; indoor sources; coagulation; deposition; dynamic aerosol model

INTRODUCTION

Ultrafine particles have been associated with adverse human health effects such as pulmonary and cardiovascular diseases (Bräuner et al. 2007; Stölzel et al. 2007). They have relatively large surface area, can carry toxic air pollutants, and have high alveolar deposition efficiency, presenting high potential to translocate to the circulation (Oberdörster et al. 2007). Given that people spend a majority of their times indoors (Klepeis et al. 2001), people are exposed to indoor UFP sources. These sources are mainly associated with occupant activities such as combustion due to cooking and using consumer products that release UFP. The UFP emissions from these indoor sources are intermittent; however, the indoor UFP concentrations due to the indoor sources are often much higher than those caused by outdoor UFP sources (Wallace and Howard-Reed 2002; Matson 2005).

Indoor UFP are subject to aerosol transport processes such as coagulation, deposition, and evaporation/condensation. In the indoor environment, ventilation is also a significant UFP loss mechanism. Several studies in the past decade have modeled indoor particle dynamics including particle emission, deposition, and penetration (Glytsos et al. 2010; Rim et al. 2010; Hussein et al. 2009^a; Hussein et al. 2009^b; Wallace 2006; Afshari et al. 2005; He et al. 2004; Thatcher et al. 2002; Lai and Nazaroff 2000; Long et al. 2001), even though only few studies have extensively monitored dynamics of size-resolved UFP in real buildings (Rim et al. 2010; Wallace 2006). Few researchers (Glytsos et al. 2010, Wallace et al. 2008; Nazaroff and Cass 1989) have investigated particle coagulation indoors, since it is quite challenging to simultaneously track the separate effects of coagulation, deposition and ventilation. Even though considerable modeling efforts have been made previously to explore aerosol coagulation theory and mathematical modeling (Jacobson 2005; Seinfeld and Pandis 2006; Whitby and McMurry 1997), experimental measurements to explore coagulation of indoor UFP are limited compared to research on outdoor UFP. Glytsos et al. (2010) monitored UFP number concentrations due to several indoor sources and indicated that strong coagulation effects shift the particle size distribution to larger sizes during candle burning, smoking, onion frying, and hair dryer operation. Dennekamp et al. (2001) reported that coagulation seems to be important for the evolution of ultrafine particles released from both gas and electric stoves.

However, a majority of earlier studies on indoor aerosol have not been able to track size-resolved coagulation of particles much smaller than 10 nm in diameter mainly due to instrument limitations; Wallace et al. (2008) measured particles as small as 2.5 nm using nano-differential mobility analysis techniques and discovered that some indoor sources such as stovetop (but not oven) cooking with gas and electric stoves emit > 90 % of the particles in the < 10 nm range.

The effect of coagulation is more important for smaller particles (< 10 nm), given their high mobility. Besides particle size, particle coagulation is also a function of Brownian motion, van der Waals force, viscosity, and fractal aggregate dimension. Glytsos et al. (2010) showed that the Brownian coagulation kernel estimated with the Fuchs correction can describe the temporal change in number concentration with reasonable accuracy for particles larger than 10 nm. Brownian motion is random particle movement caused by collisions with surrounding gas molecules and is an important mechanism for particle coagulation (Nazaroff 2004). The Fuchs correction introduces the influence of the particle size relative to the particle mean free path and yields better predictions of coagulation kinetics for small particles in the transition and free molecular regimes (Seinfeld and Pandis 2006). Along with Brownian motion, the van der Waals force enhances the rate of coagulation due to the interaction of polarized dipoles and is particularly significant for dynamics of nano-sized particles (< 50 nm). However, viscosity tends to reduce the likelihood of collision due to the fluid resistance between approaching particles (Jacobson

2005). The non-spherical shapes of UFP such as chain aggregates of a certain fractal dimension are also important in estimating particle mobility and collision rate (Kostoglou and Konstandopoulos 2001). All these elements of a coagulation model were incorporated in our earlier study of gas and electric stoves and an electric toaster oven (Wallace et al., 2008).

However, Wallace et al. (2008) made use of a theoretical deposition model for smooth surfaces based on Lai and Nazaroff (2000). In the present study, we experimentally measured and developed a size-resolved UFP deposition curve for the case of central forced fan operating, based on more than 75 days of indoor and outdoor UFP measurements in a test house on the campus of the National Institute of Standards and Technology (NIST). The experimentally measured deposition rates are considerably higher than the Lai-Nazaroff model for smooth surfaces, and thus result in increased importance of deposition. Therefore we have re-investigated UFP dynamics due to the sources (gas and electric stove) described in Wallace et al. (2008), and have added an experimental/modeling investigation of three consumer products (candle, hair dryer, and power tools). We have examined relative effects of coagulation and deposition on particle dynamics at various concentration levels. The sources considered here represent at least three different processes generating UFP: combustion, emissions from heating elements, and emissions from electric motors. Combustion of natural gas produces two types of particles: small spherical particles and larger aggregates of primary particles (Dobbins 2007). Heating elements are thought to produce UFP consisting of clusters of atoms leaving the surface (Schrupp et al. 2011; Schmidt-Ott 1988). Electric motors produce copper particles from the action of the graphite brushes on the commutator (Szymczak et al. 2007).

For these varied sources of UFP, we have developed a particle dynamics model that predicts temporal changes in size distribution following an episodic release of high concentration UFP. Given that previous studies have explored the particle emission period and reported source strengths as emission rates (Glytsos et al. 2010; Buonanno et al. 2009; Wallace et al. 2008), we focus on the decay period to examine the differential effect of coagulation compared to deposition and ventilation. For this decay period, the model analyzes three indoor particle loss mechanisms: coagulation, deposition, and ventilation. The study also performs scaling analysis with physical parameters that influence coagulation and particle size distribution.

METHODS

Dynamics of the indoor UFP size distributions were investigated using experimental measurements and analytical modeling. Experimental measurements were conducted in a full scale single-story test house, which consists of three bedrooms, two baths, a kitchen, and a dining room and a living area. The total floor area and volume of the house are 140 m² and 340 m³, respectively. In the test house, particle

monitoring tests were conducted from January 2007 to September 2010. The floor plan and details of the house layout are described in the article previously published by Rim et al. (2010). The house was furnished with limited tables, chairs, and measurement equipment, but not fully furnished with bedroom and living room furniture. During the tests, the house was occupied by one or two researchers. A summary of the experimental measurements and test conditions is provided in Table 1. For each indoor particle source, three to five tests were conducted, with variations in the particle emission period and the central forced fan operating mode.

In the tests with gas and electric stoves, particles were generated in the kitchen, the central forced fan was continuously running, and the Scanning Mobility Particle Sizer (SMPS) was in the master bedroom (MBR). All interior doors were open and the registers were open. When the central mixing fan is running, particles deposit on the air filter and duct surfaces as well as indoor surfaces. In this case, the measurement domain was the whole house. The mixing fan circulated all the air in the house through the ductwork every 10 minutes (six house volume of air per hour). We carefully examined measurement data of tracer gas (SF_6) at six different locations within the house. Using the multiple-position monitoring data, the relative standard deviation (RSD) of the SF_6 concentrations was calculated for every 10-minute interval to evaluate the uniformity of air mixing with the central mixing fan on. For the particle measurement periods, the majority of the data had RSDs $<10\%$, indicating that mixing was indeed rapid due to the central fan operation during the experiments.

For the candle, hair dryer, and power tool (belt sander and power saws) tests, particles were generated in the MBR, the door was closed, the central fan was off, mixing fans ran in the MBR and the SMPS monitored particles in the MBR. In these cases, the entire measurement area was a single room, with doors and registers sealed, so any delay due to achieving a well-mixed condition was negligible.

During the experiments, the indoor temperature ranged from $19\text{ }^{\circ}\text{C}$ to $27\text{ }^{\circ}\text{C}$ and relative humidity (RH) ranged from 17% to 60% . The air change rate was measured using an automated monitoring system in the house. A tracer gas (SF_6) was released into the living room at 4-h intervals and the decaying concentrations were monitored at 10-min intervals at six locations in the house (master bedroom, bedroom 2, living room, family room, dining room, and kitchen). The concentrations were measured using gas chromatography with electron capture detection (GC/ECD), and the air change rate was estimated by regressing the natural logarithm of the tracer gas concentration versus time. The averaged air change rates observed for the particle emission and decay period in the test house ranged from 0.14 h^{-1} to 0.51 h^{-1} . The uncertainty in the measured air change rate averaged 8% with a maximum of 15% , based on combining the two independent measurement uncertainties in quadrature: 1) the uncertainty due to incomplete air mixing plus precision (4.3% to 6.5%) and 2) the uncertainty due to calibration drift of the instrument (5%). The particle sources were turned on for 10 min to 120 min; the subsequent particle

decay was monitored for 60 minutes (twenty four 2.5-minute scans) and the dynamic analytical model was applied to the decay period.

Table 1 Test conditions

UFP Source	Test ID	Emission Period (min)	Source Location	Central Forced Fan Mode	Average indoor conditions		
					Air change rate (SD) (h ⁻¹)	Temp. (SD) (°C)	RH (SD) (%)
Gas Stove	GAS1	10	Kitchen	On	0.14 (0.05)	26.6 (1.2)	34 (2)
	GAS2	18	Kitchen	On	0.38 (0.02)	24.9 (1.6)	22 (2)
	GAS3	120	Kitchen	On	0.46 (0.08)	25.6 (1.3)	17 (1)
	GAS4	100	Kitchen	On	0.28 (0.17)	22.8 (0.4)	28 (1)
	GAS5	18	Kitchen	On	0.35 (0.10)	23.8 (1.4)	34 (1)
Electric Stove	ELEC1	10	Kitchen	On	0.26 (0.03)	21.3 (0.3)	21 (1)
	ELEC2	20	Kitchen	On	0.19 (0.01)	19.1 (0.4)	53 (5)
	ELEC3	11	Kitchen	On	0.51 (0.21)	20.6 (0.4)	23 (1)
	ELEC4	25	Kitchen	On	0.25 (0.04)	21.6 (0.6)	57 (2)
Candle	CAND1	60	MBR ^a	Off	0.18 (0.03)	24.1 (0.6)	41 (0.3)
	CAND2	60	MBR	Off	0.22 (0.01)	22.7 (0.6)	51 (4)
	CAND3	60	MBR	Off	0.20 (0.06)	24.1 (0.1)	59 (3)
Hair Dryer	HAIR1	15	MBR	Off	0.37 (0.18)	22.3 (0.5)	46 (2)
	HAIR2	15	MBR	Off	0.28 (0.09)	23.8 (1.6)	26 (2)
	HAIR3	15	MBR	Off	0.31 (0.10)	25.4 (1.9)	31 (0.9)
	HAIR4	10	MBR	Off	0.25 (0.05)	22.8 (0.6)	40 (2)
Power Tools	BELTS ^b	5	MBR	Off	0.28 (0.01)	25.6 (0.4)	12(0.1)
	PSAW ^c	2	MBR	Off	0.40 (0.02)	26.5 (0.5)	13(0.3)
	CIRSAW ^d	4	MBR	Off	0.46 (0.06)	26.6 (0.6)	13 (0.1)

a. Master Bedroom; b. Belt Sander; c. Power Saw; d. Circular Saw

The particles from each source were monitored in the master bedroom using a Scanning Mobility Particle Sizer (SMPS) (Model 3936, TSI, St. Paul, MN), which consists of an electrostatic classifier (Model 3080), a nano-DMA (Model 3085) and a water-based condensation particle counter (CPC; Model 3786). This system monitored particles ranging from 2 nm to 100 nm using two sheath flow rates of 15 L/s and 6 L/s with 10:1 sheath/aerosol flow ratio. It was extremely important to ensure that inlet sheath flow was laminar and accurately regulated, and precisely matched the exhaust flow. As small flow disturbances can result in decreased resolutions, aerosol flow rates were measured before and after each

experiment using a bubble flow rate meter. At least three measurements were taken and all three were required to be within 3 % of the desired flow rate. The sampling rate of UFPs was set at 2.5 min: 2 min of measurement and 30 s for the voltage to return to baseline. Based on the measurement efficiency curve for the water-based CPC, the minimum measurable particle size was 2 nm. In this study, a stronger radiation source (3080N electrostatic classifier, TSI, Shoreview, MN) was used to improve the charging efficiency for the smallest particles. The uncertainty of measuring particle number concentration was estimated to be 12 % based on combining the individual uncertainties of flow rate, particle charge distribution, voltage adjustment, and particle charge efficiency in quadrature (TSI 2007).

Two different sets of particle deposition rates were used depending on the fan operating mode. For the tests that were performed with the central forced fan operating, the central fan circulated all the air in the house through the ductwork every 10 minutes and air mixing was rapid. Under these circumstances, the deposition rate was a composite of at least three processes: deposition to interior surfaces in the house; removal by the furnace filter (low efficiency mechanical filter); and deposition to ductwork. The integrated deposition loss was estimated using a mass balance equation. Measuring detailed effects of the three individual deposition loss mechanisms is beyond the scope of our work and is described in Howard-Reed et al. (2003). The present study used experimental measurements of particle deposition for 25 different data sets with the central mixing fan on, each covering approximately 3 days (Rim et al. 2010). The points in Figure 1a are average deposition rates of the twenty-five 3-day periods for particles ranging from 5 nm to 100 nm; error bars are standard deviations. Using these average points, we regressed the data to estimate particle deposition rates for particle sizes from 2 nm to 5 nm. The linear logarithmic relationship yielded deposition rates for the whole house with the central mixing fan on ranging from 0.28 h⁻¹ to 6.8 h⁻¹ for particle sizes from 2 nm to 100 nm (Figure 1a).

The other set of tests were run with the central fan off. For this case, the deposition rates were derived from the Lai-Nazaroff model (2000), using a friction velocity of 3 cm/s. Typical friction velocity for indoor environments ranges from 0.3 cm/s to 3.0 cm/s, based on experimental and modeling studies on momentum boundary layer thickness and free stream bulk flows in rooms (Lai and Nazaroff 2000; Zhang et al. 1995). In our experiments with central fan off, we used mixing fans to achieve uniform air mixing in the room; therefore we used the higher value (3.0 cm/s) to model the room deposition rates. We also used surface to volume ratio (1.9 m²/m³) and specific gravity of the particles of 1.0. In this case, the deposition rate (k) ranged from 0.016 h⁻¹ to 2.0 h⁻¹ for particle sizes between 2 nm and 100 nm (Figure 1b).

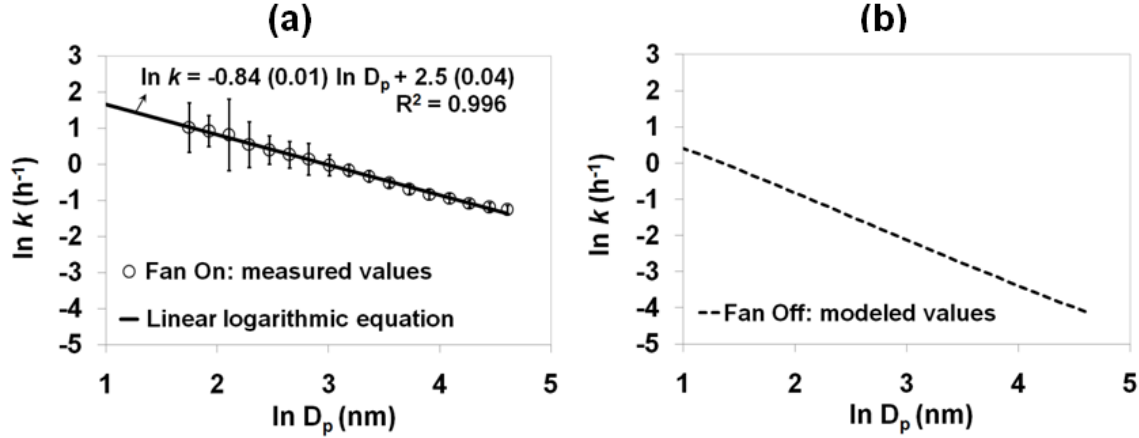


Figure 1 Double logarithmic plot of particle size (D_p) and deposition rates (k) for two different experimental conditions: (a) fan on and (b) fan off. For the fan-on condition (a), measured values (points) represent the deposition rates observed with the central fan operating. In this case, the measurement domain is the whole house and particles deposit on the air filter, duct surfaces and indoor surfaces. A linear logarithmic equation is estimated for the measurements to evaluate the deposition rates for unmeasured particle sizes. For the fan off condition (b), modeled values were derived from the Lai-Nazaroff model. In this case, the entire domain is a single room and the deposition mainly occurred at the room surfaces.

For each of the indoor sources, the source strength was calculated using particle losses due to deposition rate and air change rate, the peak total concentration, and the involved volume (Wallace et al. 2008), as shown in Equation (1)

$$S = (V(a + k) \times C_{peak}) / [1 - \exp \{-(a + k)t\}] \quad \text{Equation (1)}$$

where S = source strength [min^{-1}], V = involved volume (the whole house or MBR) [cm^3], C_{peak} = peak total concentration [cm^{-3}], a = air change rate [min^{-1}], k = total deposition rate [min^{-1}], t = particle emission period [min].

Analytical Model

The measurements of particle dynamic behavior were analyzed with an indoor aerosol model that considered coagulation, deposition, and ventilation in a single, well-mixed air volume. Effects of condensation and evaporation were not considered, even though particle condensation may have caused increases in particle mass in some tests. To address these effects, conservation of particle mass was checked and data affected by condensation and evaporation was excluded from the analysis. The governing equation for the analytical model is as follows.

$$\frac{dn(v, t)}{dt} = \frac{1}{2} \int_0^v \beta_{v-\bar{v}, \bar{v}} n(v - \bar{v}, t) n(\bar{v}, t) d\bar{v} - n(v, t) \int_0^\infty \beta_{v, \bar{v}} n(\bar{v}, t) d\bar{v} - (k + a)n(v, t)$$

Equation (2)

where $n(v, t)[m^{-3}]$ is the particle number concentration for particle volumes $[m^3]$ between v and $v + dv$ at time t [s]; $v - \bar{v}$ and \bar{v} are volumes of two coagulating particles; v is the volume of newly coagulated particle; β is the coagulation kernel (collision kernel); a is the air change rate $[h^{-1}]$; and k is the particle deposition rate $[h^{-1}]$.

The first two terms on the right hand side in Equation (2) represent the particle gain and loss due to coagulation in which two particles collide and stick together. The third term represents the particle loss due to deposition and ventilation. The effects of outdoor particle penetration were not considered for the measurement period since the peak UFP concentrations due to indoor sources are at least 30 times larger than that due to infiltrated outdoor particles.

In the present study, the concentration due to coagulation was solved by using the semi-implicit procedure in Jacobson (2005) with a volume fraction factor (f) in order to preserve volume concentrations of the particles. A volume concentration ($\rho_{q,t}$) for size bin q can be written as

$$\rho_{q,t} = \frac{\rho_{q,t-\Delta t} + \Delta t \sum_{j=1}^q \sum_{i=1}^{q-1} f_{i,j,q} \beta_{i,j} \rho_{i,t} n_{j,t-\Delta t}}{1 + \Delta t \sum_{j=1}^{N_B} (1 - f_{k,j,q}) \beta_{q,j} n_{j,t-\Delta t}} - (a + k) \Delta t \rho_{q,t-\Delta t}$$

Equation (3)

where Δt is a time step, N_B is total number of size bins, and $f_{i,j,q}$ is the volume fraction factor defined as

$$f_{i,j,q} = \begin{cases} \frac{v_{q+1} - V_{i,j}}{v_{q+1} - v_q} \frac{v_q}{V_{i,j}}, & v_q \leq V_{i,j} \leq v_{q+1}, \quad q < N_B \\ 1 - f_{i,j,q-1}, & v_{q-1} \leq V_{i,j} \leq v_q, \quad q > 1 \\ 1, & V_{i,j} \geq v_q, \quad q = N_B \\ 0, & otherwise \end{cases}$$

Equation (4)

where $V_{i,j} = v_i + v_j$ is an intermediate volume between particle size bins i and j .

The particle size distribution at the start of the decay (peak total number concentration) was used as the initial concentration profile in the model. The coagulation kernel (β) can be affected by several factors: Brownian motion, van der Waals force, viscous force, fractal geometry, convective Brownian enhancement, gravitational sedimentation, and turbulent motions. The details for calculating the coagulation kernel are described in Jacobson (2005). In the present study, we tested three coagulation models 1) Brownian motion with Fuchs correction [Brownian]; 2) Brownian motion with Fuchs correction + van der Waals and viscous forces [Brownian+VDW]; and 3) Brownian motion with Fuchs correction + van der Waals and viscous forces + fractal geometry [Brownian+VDW+Fractal24]. With

regard to fractal geometry, it was assumed that particles > 24 nm could be particle aggregates that obey a power law: $V \sim L^x$, where V is the particle volume, L is a characteristic length of the aggregate and x is the fractal dimension. Treatment of aggregates as fractal particles increases the rate of coagulation. Fractal dimensions of soot aggregates tend to range from 1.5 to 3.0 (Xiong and Friedlander 2001) and for the present study a value of 1.7 was chosen for the fractal dimension, based on a soot-aggregate study of atmospheric aerosol (Jacobson 2004).

For each of the three models considered herein, the study considered three values of the Hamaker constant: $20 k_B T$, $100 k_B T$, and $200 k_B T$ (where k_B is the Boltzmann constant and T is absolute temperature). The Hamaker constant is an empirical value that represents potential binding interaction between two colliding particles; in general it ranges from 6×10^{-20} J to 150×10^{-20} J and is greater for metal particles than organic particles (Hinds 1999). However, actual values are not known with precision.

Figures 2a – 2d illustrate the calculated coagulation kernel of any two particles between 2 nm and 100 nm. In the graph, the x- and y-axis represent the diameters of the particles. The contour values represent the collision rate between the two particle sizes. The figures indicate that the collision rate is largest for a small particle colliding with a large particle. The collision rate between two small particles, particularly < 50 nm, becomes larger when the effects of van der Waals and viscosity forces are considered (Figure 2b). The Hamaker constant also influences the coagulation kernel. Given that the Hamaker constant is a material-dependent factor that represents potential interaction between two particles, the collision rate increases as the Hamaker constant increases from $20 k_B T$ to $200 k_B T$ (Figure 2c). Fractal agglomerates lead to an increase in coagulation rate, especially for particles larger than 24 nm (Figure 2d). When a particle aggregate (> 24 nm) collides with a small particle, the rate increases up to a factor of 7 compared to the case without the fractal effect.

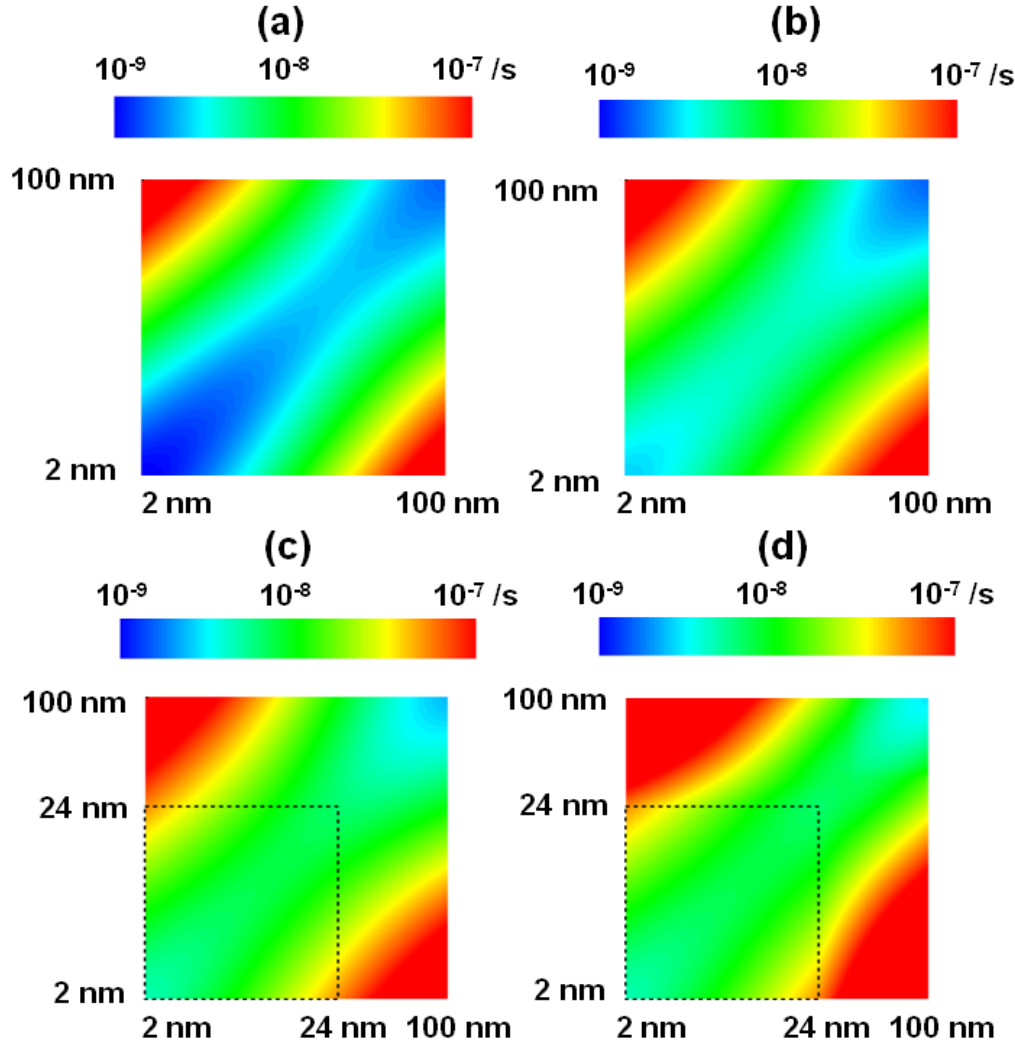


Figure 2. Coagulation kernel: (a) Brownian; (b) Brownian + VDW with Hamaker constant of $20 k_B T$; (c) Brownian + VDW Hamaker constant of $200 k_B T$; (d) Brownian + van der Waals + Fractal24 with Hamaker constant of $200 k_B T$. Adding the fractal effect does not affect the collision rate within 24 nm limit. Note that Hamaker constant affects the coagulation kernel only when van der Waals and viscosity forces (VDW) are considered.

Using the calculated coagulation kernel and the initial particle size distribution observed during the decay period, the model predicted the evolution of the particle size distribution over the subsequent period. An analytical model fitting was performed against the experimental particle decay data. The fitting error was estimated as the root-mean-square (RMS) difference in particle number between the measurement and model prediction for the eight consecutive scans (20 min) following the peak. Based on the fitting error, the present study found the best coagulation model and the optimal values of Hamaker constant that most closely match the observed evolution of the size distributions resulting from the four indoor sources. The model results also provided differential effects of coagulation and deposition on ultrafine particle size dynamics under two different cases of air mixing.

RESULTS AND DISCUSSION

Figures 3a-3f illustrate the model prediction and observations of the evolution over time of the particle size distribution due to coagulation, deposition, and ventilation for the six indoor sources studied. The particle peak occurs between 4 nm and 10 nm for the gas stove, electric stove, candle, and belt sander, while the peak is larger than 10 nm for the circular saw and hair dryer. As seen in the figures, the particle size distribution shifts over time toward larger diameters. This shift mainly results from the coagulation effect: loss of small particles and gain of large particles due to particle collisions. The coagulation effect seems to be greater for small particles than for larger particles, given the faster decrease in number concentration of small particles. The fast disappearance of small particles occurs because of the higher number concentrations and higher mobility of small particles than larger particles.

Figures 3a-3f suggest that the temporal change in particle concentration is larger when the particle concentration is higher. This result supports the second-order dependence of coagulation rate on the particle number concentration as shown in Equation 2 (Nazaroff 2004). The particle number concentration becomes lower with time and the coagulation effect decreases. For instance, with the central fan operating in the house, when the UFP number concentration (2 nm to 100 nm) decreases down to $20,000 \text{ cm}^{-3}$, the coagulation loss was approximately 20 % of the deposition loss.

Due to the greater coagulation effect with higher number concentration, the model prediction of particle size distribution shift is better with higher particle concentrations and becomes less accurate as the number concentration reduces. This decrease in accuracy with lower number concentrations is likely due to the increased effect of deposition and ventilation as well as the errors caused by local air flow effect. Also, the error accumulation with time marching in the model may also result in the decrease in accuracy with lower concentrations. However, the figures suggest that the analytical model can predict the general pattern of change in UFP size distribution, i.e., decreasing particle number and increasing geometric mean diameter with time.

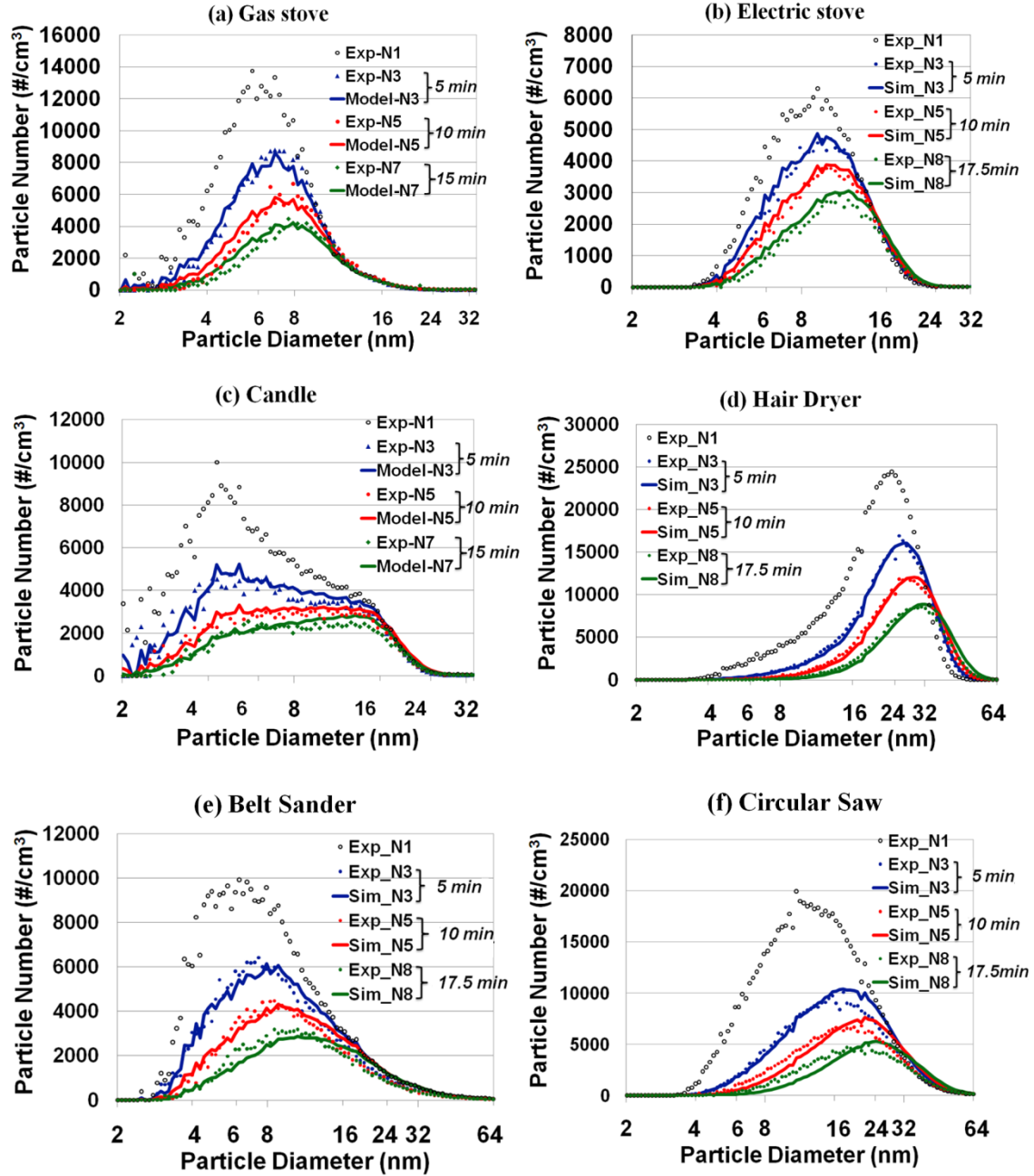


Figure 3. Examples of the analytical model fittings: (a) gas stove (GAS1), (b) electric stove (ELEC4), (c) candle burning (CAND1), (d) hair dryer (HAIR4), (e) belt sander (BELTS), (f) circular saw (CIRSAW). The numerals (N1, N3...) refer to the scans following the beginning of the decay period. The discrepancies between measurement and modeling are summarized as average root mean square errors in Table 2.

Figures 4a and 4b illustrate differential effects of coagulation, deposition, and ventilation on the temporal change in particle number concentration between successive scans following the peak, with the simulation results from the analytical model shown by the lines. The figures also show the integrated

change in total particle number observed in the experimental measurements (dots). The measurement data of the change in total particle number are scattered around the analytical model results. Figure 4a illustrates a test with the gas stove (GAS1) with the central forced fan operating. Because of the high deposition rate, the particle deposition loss is much higher than ventilation loss and nearly comparable to the coagulation loss. According to Figure 4a, the contributions of coagulation, deposition, and ventilation to the total particle number loss over the initial 2.5 min following the peak are 47 %, 48 %, and 5 %, respectively. This result suggests that when a central fan mechanically recirculates air in a building, the influence of deposition loss on the change in particle size distribution can be substantial. However, as particle concentrations decrease with time, deposition tends to dominate over coagulation. Figure 4b shows the results with the central fan turned off for one of the tests with the hair dryer (HAIR4). In this case, deposition loss only occurs on the room surfaces; therefore the deposition loss is much smaller than the coagulation loss. In this case, the contribution of coagulation, deposition, and ventilation to the total particle loss over the initial 2.5 min are 93 %, 4 %, and 3%, respectively. This result demonstrates that when the central forced air fan is off, coagulation is dominant over deposition and ventilation.

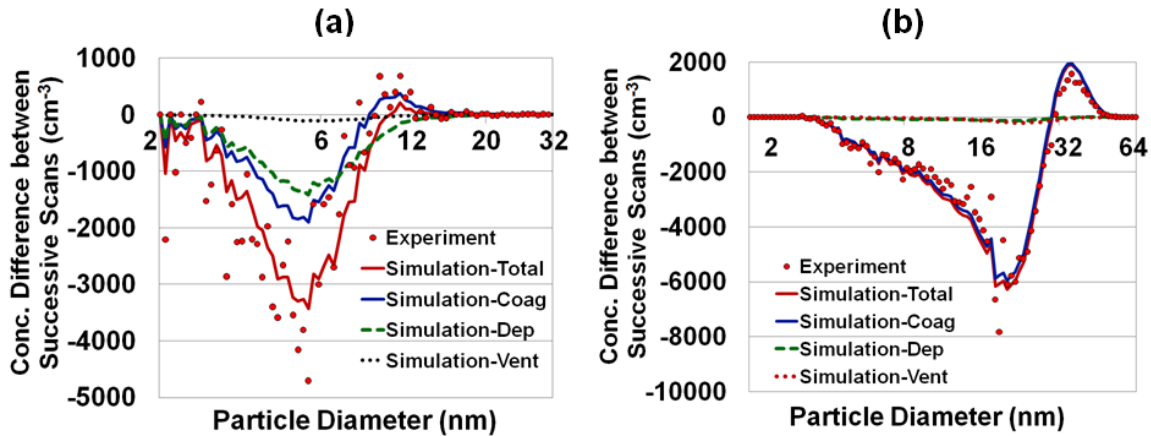


Figure 4. Size-resolved change in number concentration between first and second scan due to coagulation, deposition, and ventilation: (a) Gas stove (GAS1); (b) Hair Dryer (HAIR4). The central forced fan was operating in the GAS1 test, while it was off in the HAIR4 test. Note that the peak concentration for HAIR4 test was approximately twice that for the GAS1 test.

Table 2 summarizes the UFP size distribution and the best coagulation model/Hamaker constant for the four tested indoor sources. Based on the column showing best coagulation model, it appears that considering Brownian motion with Fuchs correction, van der Waals, and viscous forces are important in modeling indoor particle coagulation. The importance of Brownian forces and Fuchs correction for modeling particle coagulation has been shown by previous studies (Nazaroff and Cass 1989; Seinfeld and Pandis 2006; Glytsos et al. 2010). The present study results suggest that van der Waals and viscous forces

are also important for modeling indoor UFP coagulation. Van der Waals forces enhance particle coagulation due to random charge fluctuations and dipole-dipole attraction among uncharged particles. Previous studies indicated that van der Waals forces significantly enhance the rate of coagulation of particles smaller than 50 nm, while viscosity forces reduce the van der Waals enhancement due to the fluid resistance between approaching particles (Jacobson 2005; Seinfeld and Pandis 2006). The enhancement due to van der Waals forces is a function of the Hamaker constant, which varies with the particle composition.

The results in Table 2 show the best fitting values of the Hamaker constant for the indoor sources, demonstrating that the optimal value of the Hamaker constant varies with indoor source and particle composition. In general, for the gas stove and candle, a value between 20 $k_B T$ and 100 $k_B T$ produces optimal fit to the experimental data, while a constant between 100 $k_B T$ or 200 $k_B T$ seems to be more appropriate for electric stove and hair dryer. This result suggests that particles from gas stove and candle likely consist of organic matter and that those from electric stove and hair dryer are metal particles. Further studies of particle collection and imaging using Transmission Electron Microscopy (TEM) are planned to investigate more details of particle composition and binding interaction.

With regard to fractal effect, considering fractal geometry (24 nm) beside Brownian and van der Waals + viscosity forces does not substantially improve the model prediction for gas stove, electric stove, and candle. These results may be due to the fact that for those sources the geometric mean occurs at a particle size (< 10 nm), much smaller than the fractal cutoff size (24 nm). On the other hand, the fractal effect seems to be notable for some of the hair dryer tests (HAIR1 and HAIR2) and power tool test (CIRSAW), in which the peak concentration occurs at particle sizes larger than 10 nm and 22 nm. It might be possible that fractal cutoff size close to the mode of the particle distribution may result in increased effect of fractal dimension on the analytical model fitting. Nonetheless, the fractal effect is not clear for the rest of the hair dryer/power tool tests, suggesting that Brownian and van der Waals + viscosity forces are more crucial factors to consider than fractal effects in modeling the coagulation.

The best-fitting Hamaker constants found in this study agree with the findings of Wallace et al. (2008). Although Wallace et al. (2008) used only two alternate Hamaker constants of 20 $k_B T$ and 200 $k_B T$, both studies found relatively lower Hamaker constants (20 $k_B T$ – 100 $k_B T$) for gas combustion sources and higher values (100 $k_B T$ – 200 $k_B T$) for heating element sources. However, in comparison to Wallace et al. (2008), the present study suggests that particle deposition is a crucial process beside coagulation for indoor UFP dynamics especially when a central air mixing system is operating. Also, this study indicates that high UFP concentration occurs due to consumer products (candle, hair dryer, and power tools) and many of the released particles are smaller than 10 nm. However, the best-fitting coagulation model parameters vary with UFP source, size distribution, and particle composition.

Table 2. Summary of the analytical model results

Source Tested	Test ID	Source Strength ($\times 10^{12} \text{ min}^{-1}$)	Geometric Mean Diameter (nm)	Best coagulation model(s)	Best-fitting Hamaker Constant ($k_B T$)	Avg. (SD) Root Mean Square (RMS) Error ^a
Gas Stove	GAS1	11.8	5.7	Brownian+VDW	20	364 (78)
	GAS2	16.7	5.5	Brownian+VDW	20	833 (643)
	GAS3	1.91	7.1	Brownian+VDW	100	153 (52)
	GAS4	2.22	5.5	Brownian+VDW	20	339 (70)
	GAS5	14.0	4.8	Brownian+VDW	20	549 (141)
Electric Stove	ELEC1	3.43	8.8	Brownian+VDW	200	148 (21)
	ELEC2	0.48	4.8	Brownian+VDW	20-200 ^b	88 (15)
	ELEC3	6.18	7.4	Brownian+VDW	100 or 200	194 (127)
	ELEC4	2.64	9.1	Brownian+VDW	200	179 (15)
Candle	CAND1	0.227	4.3	Brownian+VDW	100	359 (115)
	CAND2	0.044	6.6	Brownian+VDW	20 or 100	122 (30)
	CAND3	0.137	4.5	Brownian+VDW	100	262 (49)
Hair Dryer	HAIR1	0.512	14	Brownian+VDW + Fractal24	200	166 (27)
	HAIR2	1.11	22	Brownian+VDW + Fractal24	100	289 (62)
	HAIR3	0.916	22	Brownian+VDW	100 or 200	654 (186)
	HAIR4	2.15	23	Brownian+VDW	100	426 (16)
Power Tools	BELTS	2.44	4.8	Brownian+VDW	100	346 (63)
	PSAW	13.3	8.8	Brownian+VDW	100	784 (84)
	CIRSAW	5.84	11	Brownian+VDW + Fractal24	100	679 (42)

a. The mean (SD) RMS difference in particle number between the measurement and model prediction are based on the fittings of eight consecutive particle size distributions with an interval of 2.5 minutes during the initial 20 minutes of the decay period.

b. This indicates that there was little difference ($< 5\%$) in the model results when varying the Hamaker constant from 20 $k_B T$ to 200 $k_B T$.

CONCLUSIONS

The present study investigated dynamic aerosol processes for evolution of UFP size distribution following episodic releases of UFP in a residential building. Using experimental measurements of time-varying, size-resolved particle concentrations, the temporal change in UFP size distribution was observed for five different indoor sources. Using the experimental data, the present study developed a particle dynamics model that analyzed the differential influence of coagulation, deposition, and ventilation on particle size dynamics. The study confirms previous findings that coagulation is an important aerosol process affecting UFP dynamics when particle concentrations are high due to indoor sources such as cooking (Wallace et al., 2008), and has extended these findings to several additional sources (candles, hair dryers, power tools). The coagulation effect was higher with smaller UFP and higher number concentrations. The study results suggest that deposition should be considered depending on the operating mode of the central mechanical fan. When the central fan was turned on, particle deposition loss was substantial and comparable to the coagulation loss. When it was off, deposition loss was not significant and coagulation was the primary cause for the shift of UFP size distribution. Compared to coagulation and deposition, air change was a minor effect in these tests. Therefore, in terms of aerosol modeling for episodic indoor UFP releases, coagulation should be considered during high concentration periods ($> 20,000 \text{ cm}^{-3}$), while deposition should be treated as a major loss mechanism when air recirculates through ductwork or mechanical systems.

The analytical model that considers Brownian, van der Waals/viscosity forces and fractal geometry along with deposition and ventilation can predict with reasonable accuracy the temporal particle size distribution due to the indoor sources. The study results also suggest that using an appropriate Hamaker constant and deposition rate is essential for an accurate prediction of indoor UFP dynamics.

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