# Nanowire Electromechanical Logic Switch 

Qiliang Li *, Member, IEEE, Curt A. Richter, Senior Member, IEEE, Hao D. Xiong, and John S. Suehle, Senior Member, IEEE<br>Semiconductor Electronic Division, National Institute of Standards and Technology (NIST), Gaithersburg, MD 20899, USA.


#### Abstract

We present the integration and characterization of nanowire electromechanical switches consisting of chemical-vapor-deposition grown silicon nanowires suspended over metal electrodes. The devices operate with the suspended part of the nanowire bent to touch metal electrode via electromechanical force by applying voltage. The reversible switching, high on/off current ratio, small subthreshold slope and low switching energy compared to current Si CMOS make the switches very attractive for logic device application. In addition, we have developed a physical model to simulate the switching characteristics and extract the material properties.


Keywords - Electromechanical Switch, Silicon Nanowire, Transistor, Subthreshold Slope

## I. Introduction

During the past three decades of great success, silicon Complementary Metal-Oxide-Semiconductor field effect transistor (CMOS FET) has been aggressively scaled down to achieve better performance and lower power consumption. Beyond the CMOS scaling, which becomes more difficult in the future, next-generation logic device requires alternative strategies, such as resonant tunneling devices, single electron transistors, molecular and spin devices [1-5]. However, the operation of these devices still represents charge-based mechanism with the similar thermodynamic limit with CMOS. New technologies based on something different than electronic charge, such as Nano-Electromechanical System (NEMS) and optical computing, may extend the functionalities and open up exponential opportunity for future electronic circuitry [1]. We have previously reported the fabrication of a simple NEMS test structure consisting of a Si nanowire (SiNW) suspended over metal electrodes [6]. In this work, we demonstrated the integration and characterization of 2-terminal (2-T) and 3-terminal (3-T) electromechanical switches (EMSs). These switches are turned on/off via the electromechanical force by applying a voltage between a SiNW suspended as a cantilever and bottom metal electrodes. The switching of the EMS, which can be operated as a field effect transistor, depends on the nanowire material's mechanical properties and the device geometry. The reversible switching, small subthreshold slope, low switching energy and large on/off ratio indicated that the SiNW EMS is very attractive for logic application. In addition, we have developed a model to investigate the physics of the electromechanical switching and mechanical properties of nanowires. It should be noted that nanowire NEMS has been used to investigate the mechanical properties of nanowires and nanotubes [7, 8], and

NEMS switch has been fabricated with carbon nanotubes by using E-Beam lithography [9]. A suspended SiNW switching system used in this study may provide an alternative and perhaps simpler strategy for achieving electromechanical switching and characterizing the material properties of chemically-grown nanowires.

## II. Fabrication of Electromechanical Switches

The SiNWs used in this study (Figure 1a) were prepared by using low-pressure chemical vapor deposition (LPCVD) under $500 \mathrm{mTorr} \mathrm{SiH}_{4}$, at $450^{\circ} \mathrm{C}$, with thin Au films ( $\sim 2 \mathrm{~nm}$ ) as the catalyst via a vapor-liquid-solid mechanism [10, 11]. Nanowires of 20 nm to 300 nm in diameter and $2 \mu \mathrm{~m}$ to 150 $\mu \mathrm{m}$ in length were obtained. From transmission electron microscopy (TEM) characterization (data not shown), the SiNWs with a hexagonal cross-section grow along the expected [111] direction [12] with a spacing between the (111) planes, which are oriented perpendicular to the SiNW growth direction $\approx 0.31 \mathrm{~nm}$ as determined from TEM images.


Figure 1 (a) SEM image of SiNWs grown from patterned Au catalyst on $\mathrm{SiO}_{2}$. (b) SiNWs were removed into a suspension of DI water. (c) The SiNWs were placed on template substrate and dried, and then manipulated by SNMS.

The nanowire cantilever switches were formed by manipulating SiNWs to align them onto a device substrate by using a single nanowire manipulation system (SNMS) [6] and compatible fabrication processes. First, the SiNWs were removed into a suspension of DI water by sonication ( 1 min ) (Fig. 1b). The SiNW solution was dispersed onto a template substrate and dried for 4 hours under vacuum at 100 mTorr (Fig. 1c).

As shown in the schematic drawing (Fig. 2a and 2b), the SiNWs were then manipulated and placed in pre-defined locations on a device substrate by using a lab-made SNMS so

[^0]that one side of the SiNWs lie on a thicker oxide ( $\sim 600 \mathrm{~nm}$ in thickness) and the other side is suspended over metal electrode


Figure 2 (a) A SiNW was picked up and transferred to the pre-defined location of a device substrate by the single nanowire manipulation system probes. (b) The SiNW was patterned with metal contacts (e.g., 3-T EMS has source, drain and gate contacts.) by compatible photolithographic processes.
(for 2-T EMS) or a thinner oxide ( $\sim 200 \mathrm{~nm}$ in thickness) (for 3-T EMS). Two kinds of switches (2-T and 3-T switches) were prepared in this study by using slightly different processes. For the fabrication of 2-T EMSs, the suspended gap was obtained by dry-etching 400 nm of the $600-\mathrm{nm}$ thermal $\mathrm{SiO}_{2}$ followed by $\mathrm{Au} / \mathrm{Al}(10 \mathrm{~nm} / 10 \mathrm{~nm})$ deposition and lift-off to form the bottom electrode of the switch. After the alignment of a SiNW ( 110 nm in diameter) with SNMS, a second-level metal ( 20 nm / 60 nm of $\mathrm{Au} / \mathrm{Al}$ ) was patterned on the part of the SiNW lying on the $600-\mathrm{nm}$ oxide to clamp the nanowire in place and form the top electrode. For the fabrication of 3-T EMSs, the suspending gap was obtained by dry-etching 400 nm of the $600-\mathrm{nm}$ thermal $\mathrm{SiO}_{2}$. After the manipulation and alignment of SiNWs to the designed locations, the bottom electrodes (drain and gate) and top electrode (source) were formed in a single step (see Fig. 2b). All the metal patterns were obtained by using regular photolithographic and metal lift-off processes with an Au layer on the top of the Al layer.

## III. Characterization of Electromechanical SWITCHES

In this work, we have prepared 2-T and 3-T EMSs acting similarly as a current switch and transistor, respectively. A schematic drawing and SEM images of a $2-\mathrm{T}$ switch are shown in Figure 3 (a) and (b), respectively. The bottom and top electrodes are assigned as the drain and source of the 2-T EMS as shown in the images. The typical switching currentvoltage characteristics of a 2-T EMS are shown in Figure 3 (c) with an inset schematically showing a switch On/Off diagram. The 2-T EMS stays off until the applied voltage between the source and drain electrodes is raised above the threshold
voltage, $\mathbf{V t},(\mathrm{Vt}=3.8 \mathrm{~V}$ for the example shown in Figures 3 (c)) and a significant increase in current is observed. During


Figure 3 (a) Schematic of a 2-T EMS; (b) SEM image of a 2-T EMS; (c) the typical switching current-voltage characteristics of the 2-T EMS with an inset of switch On/Off diagram.
the operation, the suspended part of SiNW was bent down to touch the bottom electrode by the electromechanical force as it reached a snap-down point at Vt. The On/Off ratio of current measured at 9 V and 3.8 V is about $10^{3}$. Such reversible current-voltage switching is observed at both negative and positive applied voltages.

This 2-T EMS is a very simple structure and an efficient switch; however, it has two disadvantages which limit its applications. First, the on-state current is relatively low due to the high contact resistance between the intrinsically doped semiconducting SiNW and the bottom metal electrode. Second, it cannot be operated as a gate-controlled transistor with gain. To solve these problems, we have fabricated a 3-T EMS with self-aligned source, drain and gate electrodes. The schematic of a 3-T EMS is shown in Figure 4 (a). After the manipulation and alignment of SiNWs to the designed locations, the bottom electrodes (drain and gate) and top electrode (source) were formed in a single step. This differs from the fabrication of 2T EMSs which have the bottom electrode formed before the SiNW alignment. The same types of metals have been
deposited for both the 2-T EMSs and 3-T EMSs. However, during the fabrication of the 3-T EMS electrodes, the metals


Figure 4 (a) Schematic of a 3-T EMS; (b) SEM image of two 3-T EMSs; (c) the typical switching current-voltage characteristics of the 3-T EMS in the right side of the SEM image. The top inset shows the switch On/Off diagram and the bottom inset shows the hysteresis of $\mathrm{I}_{\mathrm{DS}}-\mathrm{V}_{\mathrm{GS}}$ as the gate voltage is swept from $\mathrm{V}_{\mathrm{GS}}=0 \mathrm{~V}$ to 9.0 V and then back to 0 V at $\mathrm{V}_{\mathrm{DS}}=2.0 \mathrm{~V}$.
are also deposited on the suspended nanowires. This metal directly deposited on the SiNW is expected to lower (i.e., improve) the contact resistance between the SiNW and the bottom drain electrode when the switch is closed. The two bottom electrodes and the top electrode were assigned as the drain, gate and source of the switch as shown in Figure 4 (a). During the operation of a 3-T EMS, the gate voltage is controlled independently of the drain voltage, which is similar to a conventional MOSFET.

A SEM image of a 3-T EMS is shown in Figure 4 (b) with the assigned electrodes. Figure 4 (c) shows typical switching current-voltage characteristics of a 3-T EMS with the top inset schematically illustrating the On/Off configurations. As the gate to source voltage $\left(\mathbf{V}_{\mathbf{G S}}\right)$ increases with a constant drain to source voltage ( $\mathbf{V}_{\mathbf{D S}}$ ), the electrostatic force applied to the SiNW increases and bends the wire down to the bottom electrodes. As the suspended end of the SiNW touches the drain electrode at a threshold voltage (e.g., $\mathrm{Vt}=7.9 \mathrm{~V}$ at $\mathrm{V}_{\mathrm{DS}}$ $=1.0 \mathrm{~V}$ for the device shown in Figure 4 (c)), the drain-source current ( $\mathrm{I}_{\mathrm{DS}}$ ) is turned on. Since the gate electrode is closer than the drain electrode to the source electrode in the set-up, the SiNW bends over the gate electrode with it initially touching the drain to close the switch. As shown in the bottom inset of Figure 4 (c), an $\mathrm{I}_{\mathrm{DS}}-\mathrm{V}_{\mathrm{DS}}$ hysteresis is observed when the gate voltage is scanned from $\mathrm{V}_{\mathrm{GS}}=0 \mathrm{~V}$ to 9.0 V (which above Vt ) then back to 0 V . This hysteresis arises from the surface electrostatic force between the SiNW and the drain electrode surface. In the 3-T configuration, Vt is determined by the electrostatics between the nanowire and both the gate and drain electrodes. A higher $\mathrm{V}_{\mathrm{DS}}$ leads to a smaller Vt due to the larger total attractive force associated with the drain electrode in combination with the gate electrode. Compared to the Vt of a 2-T EMS, a much higher $\mathrm{V}_{\mathrm{GS}}$ is needed in order to pull the SiNW down in contact with the gate electrode in a 3-T EMS due to the much smaller gate electrode area. The 3-T EMS has an On/Off current gain of $\sim 10^{4}$ and a subthreshold slope of $\sim 75 \mathrm{mV} / \mathrm{dec}$ at a $\mathrm{V}_{\mathrm{DS}}=1.0 \mathrm{~V}$ if it is considered as a conventional field effect transistor. The improved lower contact resistance between the metal-coated SiNW and drain of the 3-T EMS leads to a higher on-current and the larger observed On/Off ratio when compared with 2-T EMSs. The subthreshold slope and current On/Off gain can be further improved by replacing the SiNWs with more conductive nanowires.


Figure 5 Schematic of a suspending SiNW over metal electrode with the suspended height $\boldsymbol{h}$, length $\boldsymbol{L}$, radius of bending curvature $\boldsymbol{\rho}$, and the charge of unit area at $(\mathrm{x}, 0, \mathrm{z})$ on the bottom electrode applying a force on the SiNW at $(0$, $\mathrm{h}, l$ ). Bottom right corner shows the diagram of a hexagonal cross-section.

## IV. Physical Model of Electromechanical Switches

The threshold voltage for these EMS to "close" the switch and turn on the current is determined by the exact geometry of the SiNW and its mechanical properties (e.g., Young's modulus: $\boldsymbol{E}$ ). When the maximum moment due to electrostatic force $\left(\boldsymbol{M}_{e}\right)$ is greater than the bending moment $\left(\boldsymbol{M}_{\boldsymbol{m}}\right)$ of the SiNW, the SiNW will be bent down and in contact with the bottom electrode: the EMS will be closed. A simple estimate of Vt can be made by assuming the initial "snap-down" of the SiNW occurs when $\boldsymbol{M}_{\boldsymbol{e}}=\boldsymbol{M}_{\boldsymbol{m}}$. We illustrate this estimation here for the example of a SiNW with a cylindrical geometry with diameter $(\boldsymbol{R})$, suspending length $(\boldsymbol{L})$, and suspending height ( $\boldsymbol{h}$ ). These parameters are shown in the schematic diagram of a suspending SiNW over metal electrode (Fig. 5). The potential of the system is [13]
$\psi(x, y)=\lambda / 4 \pi \varepsilon_{0} \cdot \ln \left[x^{2}+(y+h)^{2} / x^{2}+(y-h)^{2}\right]$
where $\lambda$ is the charge density per unit length of the suspended $\operatorname{SiNW}, \varepsilon_{0}$ is the permittivity of vacuum, $\mathrm{x}, \mathrm{y}$ and z are the coordinates shown in Figure 5.

The surface charge density $\boldsymbol{\sigma}(\boldsymbol{x})$ of bottom electrode is therefore:
$\sigma(x)=-\left[\frac{\partial \psi}{\partial y}\right]_{y=0} \cdot \varepsilon_{0}=C \cdot V_{t} \cdot h / \pi L\left(x^{2}+h^{2}\right)$
The electrostatic force between $\partial l$ of the Si nanowire and $\partial x \cdot \partial z$ of the bottom electrode can be expressed as:
$\partial F=Q_{1} \cdot Q_{2} / 4 \pi \varepsilon_{0} r^{2}$,
where $\boldsymbol{Q}_{1}$ and $\boldsymbol{Q}_{2}$ are the charges in the differential bottom electrode surface $\partial x \cdot \partial z$ and the charges in $\partial l$ of the Si nanowire at $(0, h, l)$; and $\boldsymbol{r}$ is the distance from the bottom surface point $(x, 0, z)$ to the Si nanowire at point $(0, h, l)$ is: $r^{2}=x^{2}+h^{2}+(z-l)^{2}$. The direction of $\partial F$ is at the $\boldsymbol{r}$ direction. We have

$$
\begin{aligned}
& Q_{1}=\sigma(x) \cdot \partial x \cdot \partial z \\
& Q_{2}=C \cdot V_{t} \cdot \partial l / L
\end{aligned}
$$

The electrostatic force applied on the Si nanowire at $(0, h, l)$ by the charges in the bottom electrode surface $\partial x \cdot \partial z$ at the perpendicular direction $\left(\partial F_{\perp}\right)$ is therefore:
$\partial F_{\perp}=\partial F \cdot h / r=\sigma(x) C \cdot V_{t} \cdot h \cdot \partial x \cdot \partial z \cdot \partial l / 4 \pi \varepsilon_{0} r^{3} \cdot L$
So the electrostatic moment on the Si nanowire at point $(0, h, l)$ due to $\partial F_{\perp}$ is
$\partial M_{e}=l \cdot \partial F_{\perp}=\sigma(x) \cdot C \cdot V_{t} \cdot h \cdot l \cdot \partial x \cdot \partial z \cdot \partial l / 4 \pi \varepsilon_{0} r^{3}$
The maximum moment due to electrostatic force $\left(\boldsymbol{M}_{\boldsymbol{e}}\right)$ is therefore:
$M_{e} \approx\left(C^{2} V_{t}^{2} h / 4 \pi^{2} \varepsilon_{0} L^{2}\right) \cdot I n$
where the unitless integration term $\operatorname{In}$ is

In $=\int_{a}^{b} \int_{0}^{L} \int_{0}^{L} \partial x \partial z \partial l \cdot h \cdot l /\left(x^{2}+h^{2}\right)\left[x^{2}+h^{2}+(z-l)^{2}\right]^{3 / 2}$
$a$ and $b$ are $-10 \mu \mathrm{~m}$ and $15 \mu \mathrm{~m}$, respectively, representing the distance from the SiNW to the left and right edges of the bottom electrodes as shown in Fig. 3b. In is numerically calculated as 467 for this specific geometry. In decreases dramatically as $a$ and $b$ increase to $1 \mu \mathrm{~m}$. For a SiNW with a cylindrical cross-section suspending over metal surface the capacitance, $\boldsymbol{C}$, of the system is [13]:

$$
\begin{equation*}
C \approx 2 \pi \varepsilon_{0} L / \cosh ^{-1}(2 h / R) \tag{7}
\end{equation*}
$$

which is a slight over-estimation for a SiNW over a semifinite or large finite metal surface. For the geometry here, $\boldsymbol{C}$ is calculated as $\approx 2.0 \times 10^{-16} \mathrm{~F}$. A more detailed theoretical study of $\boldsymbol{C}$ for a similar suspended nanowire system has been reported [14], and more precise calculations of $\boldsymbol{C}$ based on special nanowire geometries [15] and the electric field effect [16] have been reported.
With $M_{m}=E \cdot I / \rho$ [17], the Young's modulus $\boldsymbol{E}, V \boldsymbol{t}$ and switching energy ( $\Delta \mathbf{W}_{\mathbf{S}}$ ) are expressed as following:
$E \approx\left(C^{2} V_{t}^{2} h \rho / 4 \pi^{2} \varepsilon_{0} L^{2} I\right) \cdot$ In
$V_{t} \approx \sqrt{4 \pi^{2} \varepsilon_{0} L^{2} E I / I n \cdot C^{2} h \rho}$
$\Delta W_{S}=\frac{1}{2} C V_{t}^{2}=2 \pi^{2} \varepsilon_{0} L^{2} E I / C h \rho \cdot$ In
where $\boldsymbol{I}$ is the moment of inertia of the SiNW, $\rho$ is the radius of curvature and $\boldsymbol{E I}$ is the flexural rigidity of the SiNW [17]. $\rho$ can be expressed as $\rho=\left(h^{2}+L^{2}\right) / 2 h$. Assuming the SiNW has a cylindrical cross-section, $I$ can be calculated as: $I=\frac{\pi}{64} \times R^{4}$. Thus, for the example shown in Figure 2, $\boldsymbol{E}$ of a cylindrical SiNW extracted by this method is $\approx 118 \mathrm{G} \mathrm{Pa}$, which is within the range of the reported values. [18]
$\boldsymbol{C}, \boldsymbol{V} \boldsymbol{t}$, and switching energy strongly depend upon the details of the bottom electrode geometry. It should be noted that the accuracy of the Young's modulus extracted by this simple estimation also depends on how well the equations represent the specific geometry of the SiNWs and the electrode. For example, for a SiNW with a hexagonal cross section (see the inset of Fig. 5, which more closely represents the SiNWs grown in $[111]$ direction [15] than the cylindrical example) with a side length, $w=\frac{1}{2} R$, the capacitance is calculated as $\approx$ $1.87 \times 10^{-16} \mathrm{~F}$ by following the reported approach [15]. The moment of inertia is $I=\frac{5 \sqrt{3}}{256} \times R^{4}$. For this example, of the SiNW with a hexagonal cross-section, the $\boldsymbol{E}$ (proportional to $\left.C^{2} / I\right)$ is calculated as 150 G Pa . Thus, the estimated Young's
modulus is different by $\approx 24 \%$ for these two different geometries. In addition, the electric field will also affect the accuracy of the extracted Young's modulus. [16] The switching energy of the 2-T EMS shown in Figure 2 is about $2.0 \times 10^{-14} \mathrm{~J}$. Theoretically, with $\boldsymbol{E}=107 \mathrm{GPa}$ (for the bulk Si), $\mathrm{L}=1 \mu \mathrm{~m}, \mathrm{R}=10 \mathrm{~nm}$ and $\mathrm{h}=10 \mathrm{~nm}$, the switching energy is as low as $3.7 \times 10^{-19} \mathrm{~J}$. In addition, based on the developed model and equation (9), ideally the device operation is unaffected by the temperature and the subthreshold slope has no limit. While there are many subtleties associated with using this technique to derive the accurate value of the Young's modulus, it can be used to precisely compare relative changes in E for nanowires with different structural properties (such as a change in radius for nanowires of a known material or a comparison of nanowires with the same diameter formed from different materials).

## V. CONCLUSIONS

In summary, we have fabricated 2-T and 3-T electromechanical switches consisting of suspended nanowires over metal electrodes. We have shown that the EMS devices are attractive for the following reasons. (1) The devices have no power consumption in the off-state and very small switching energy ( $\sim 10^{-14} \mathrm{~J}$ ). Virtually no current passes through the switch until the SiNW is mechanically "bent" and brought into contact with the bottom electrode and a large oncurrent is observed. (2) The devices have small subthreshold slope ( $\sim 75 \mathrm{mv} / \mathrm{dec}$ ). Theoretically, the switch with more conductive nanowire (e.g., metal nanowires) will have much smaller subthreshold slope than that of CMOS FET (i.e. $\sim 60$ $\mathrm{mv} / \mathrm{dec}$ ) [19]. (3) The switches are turned on/off reversibly with large On/Off current gain $\left(\sim 10^{4}\right)$.

## Acknowledgment

Contribution of the National Institute of Standards and Technology is not subject to U.S. copyright. The authors are grateful to the support of the NIST Office of Microelectronics Programs and NIST Semiconductor Electronics Division.

## REFERENCES

[1] ITRS, Emerging Research Devices, 2005 edition.
[2] R. H. Chen, A. N. Korotkov, and K. K. Likharev, "Single-electron transistor logic," Applied Physics Letters, vol. 68, pp. 1954-1956, 1996.
P. Fay, L. Jiang, Y. Xu, G. H. Bernstein, D. H. Chow, J. N. Schulman, H. L. Dunlap, and H. J. De Los Santos, "Fabrication of monolithically-integrated InAlAs/InGaAs/InP HEMTs and InAs/ $\mathrm{AlSb} / \mathrm{GaSb}$ resonant interband tunneling diodes," Ieee Transactions on Electron Devices, vol. 48, pp. 1282-1284, 2001.
D. E. Nikonov and G. I. Bourianoff, "Spin gain transistor in ferromagnetic semiconductors - The semiconductor Blochequations approach," Ieee Transactions on Nanotechnology, vol. 4, pp. 206-214, 2005.
[5] J. M. Seminario, P. A. Derosa, L. E. Cordova, and B. H. Bozard, "A molecular device operating at terahertz frequencies: Theoretical simulations," Ieee Transactions on Nanotechnology, vol. 3, pp. 215-218, 2004.
[6] Q. Li, S.-M. Koo, C. A. Richter, M. D. Edelstein, J. E. Bonevich, J. J. Kopanski, J. S. Suehle, and E. M. Vogel, "Precise Alignment of single Nanowires and Fabrication of Nanoelectromechanical Switch and Other Test Structures," Ieee Transactions on Nanotechnology, vol. 6, 2007.
[7] Z. L. Wang, Z. R. Dai, R. P. Gao, and J. L. Gole, "Measuring the Young's modulus of solid nanowires by in situ TEM," Journal of Electron Microscopy, vol. 51, pp. S79-S85, 2002. E. W. Wong, P. E. Sheehan, and C. M. Lieber, "Nanobeam mechanics: Elasticity, strength, and toughness of nanorods and nanotubes," Science, vol. 277, pp. 1971-1975, 1997.
[9] J. E. Jang, S. N. Cha, Y. Choi, G. A. J. Amaratunga, D. J. Kang, D. G. Hasko, J. E. Jung, and J. M. Kim, "Nanoelectromechanical switches with vertically aligned carbon nanotubes," Applied Physics Letters, vol. 87, pp. 163114, 2005.
[10] R. S. Wagner and W. C. Ellis, "Vapor-Liquid-Solid Mechanism of Single Crystal Growth," Applied Physics Letters, vol. 4, pp. 89-90, 1964.
[11] E. I. Givargizov, "Fundamental Aspects of Vls Growth," Journal of Crystal Growth, vol. 31, pp. 20-30, 1975.
[12] A. Morales and C. Lieber, "A laser ablation method for the synthesis of crystalline semiconductor nanowires," SCIENCE, vol. 279, pp. 208-211, 1998.
[13] P. M. Morse and H. Feshbach, Methods of Theoretical Physics, vol. 2. New York: McGraw-Hill, 1953.
[14] Z. Tang, Y. Xu, G. Li, and N. R. Aluru, "Physical models for coupled electromechanical analysis of silicon nanoelectromechanical systems," Journal of Applied Physics, vol. 97, pp. 114304, 2005.
O. Wunnicke, "Gate capacitance of back-gated nanowire fieldeffect transistors," Applied Physics Letters, vol. 89, pp. 083102, 2006.
[16] X. J. Zheng and L. L. Zhu, "Theoretical analysis of electric field effect on Young's modulus of nanowires," Applied Physics Letters, vol. 89, pp. 153110, 2006.
[17] J. M. Gere and S. P. Timoshenko, Mechanics of Materials, Third edition. BOSTON: PWS-KENT, 1991.
M. Menon, D. Srivastava, I. Ponomareva, and L. A. Chernozatonskii, "Nanomechanics of silicon nanowires," Physical Review B, vol. 70, pp. 125313, 2004.
[19] T. Taur and T. H. Ning, Fundamentals of Modern VLSI Devices. New York: Cambridge University Press, 1998.


[^0]:    * Contact author: Qiliang.Li@nist.gov

